

EXTRA S_{11} AND P_{13}

IN THE

HYPER CENTRAL CONSTITUENT

QUARK MODEL

(hCQM)

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M. G.

L. Tiator

M. De Sanctis

GE

MZ

RM

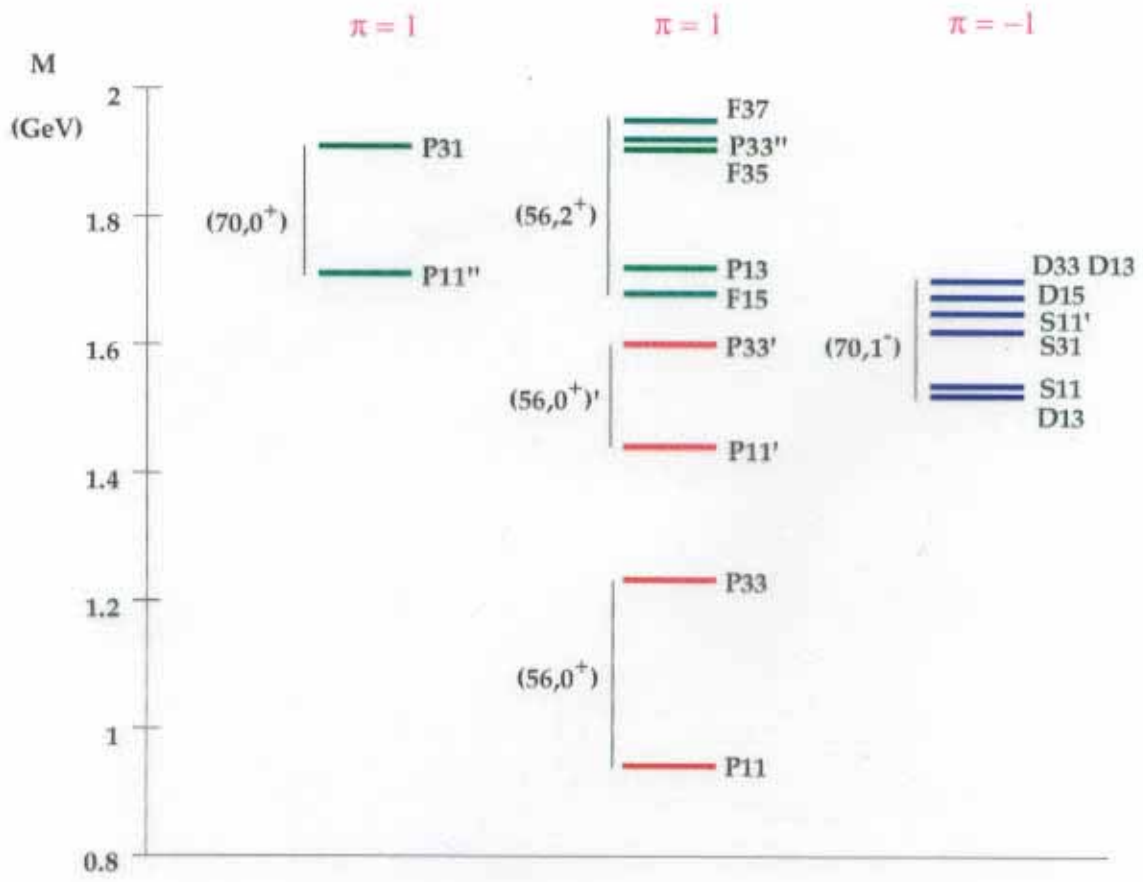
M. Giannini

NSTAR 2002

PITTSBURGH

11 Oct. 2002

4 * PDG
3 *

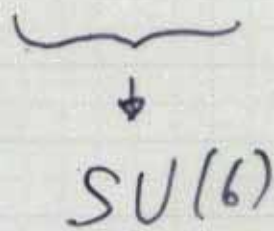


3q Wave function

$$\Psi_{3q} = \Theta_{\text{colour}} \times \chi_{\text{spin}} \times \Phi_{\text{iso}} \times \Psi_{\text{space}}$$

$$SU(3)_c \quad SU(2) \quad SU(3)_f \quad O(3)$$

Spin-flavour independent quark interaction



Θ colour singlet \rightarrow antisymm.

$SU(6)$ wf: same symmetry
 $O(3)$

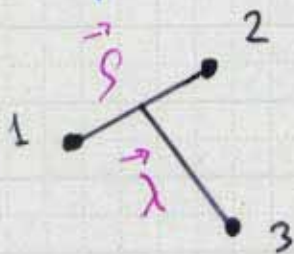
for 3 objects

- A
- MA
- MS
- S

$\chi \quad S = \frac{1}{2}, \frac{3}{2}$

$\Phi \quad T = \frac{1}{2}, \frac{3}{2}$

ψ space



Jacobi coordinates

$$\vec{s} = \frac{1}{\sqrt{2}} (\vec{r}_1 - \vec{r}_2)$$

$$\vec{\lambda} = \frac{1}{\sqrt{6}} (\vec{r}_1 + \vec{r}_2 - 2\vec{r}_3)$$

• six degrees of freedom

$$s, \Omega_s, \lambda, \Omega_\lambda$$

• new coordinates \rightarrow hyperspherical

$$x = \sqrt{s^2 + \lambda^2}$$

hyper radius ("size")

$$t = \arctg(\lambda/s)$$

hyper angle ("form")

$$+ \Omega_s, \Omega_\lambda$$

$$H_{3q} = 3m - \frac{1}{2m} (\Delta_s + \Delta_\lambda) + V(\vec{s}, \vec{\lambda})$$

$\rightarrow C_2(0(6))$

$$\frac{\partial^2}{\partial x^2} + \frac{5}{x} \frac{\partial}{\partial x} - \frac{L^2(\Omega_s, \Omega_\lambda, t)}{x^2}$$

• L^2 eigenfunctions

$$Y_{[\Omega_s, \Omega_\lambda, t]}(\Omega_s, \Omega_\lambda, t)$$

$$\times Y_{\ell_s m_s} Y_{\ell_\lambda m_\lambda} P_{\ell_t}(t)$$

Jacobi polynomials

eigenvalues

$$\gamma \cdot (\gamma + 4) \quad \gamma = 2n + \ell_s + \ell_\lambda$$

• wf basis

$$\psi(x) \cdot Y_{[\Omega_s, \Omega_\lambda, t]}(\Omega_s, \Omega_\lambda, t)$$

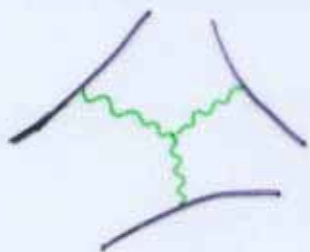
Hypercentral hypothesis

$$V(\mathbf{x})$$

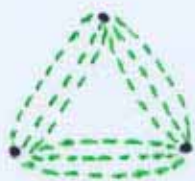
3-body potential

• motivations:

- QCD fundamental mechanisms



- flux-tube model



$$\bullet \sum_{i < j} V(\vec{r}_{ij}) \cong V_0(\mathbf{x}) + \dots$$

hypercentral approximation

$O(6)$ -invariance

"hypercentral potentials"

$$V = V(x)$$

→ factorization

$$\Psi(x, t, \Omega_p, \Omega_\lambda) = \Psi(x) Y_{[\gamma, l_1, l_2]} \\ (\text{"dynamics"}) \quad (\text{"geometry"})$$

"hypercentral equation"

$$-\frac{1}{2m} \left[\frac{d^2 \psi}{dx^2} + \frac{5}{x} \frac{d\psi}{dx} - \frac{\gamma(\gamma+4)}{x^2} \psi \right] + V(x) \psi = E \psi$$

$$\psi = \psi_{\nu \gamma}(x)$$

↓ ↗ grand angular quantum number

hyperradial quantum number
(hyperradial excitation)

$V(x)$: three-body force

two-body potential in
hypercentral approximation

$$\left\langle \sum_{i < j} (r_{ij})^n \right\rangle_{av.} \propto x^n$$

P_{11} P_{11} P_{11}

$$\begin{pmatrix} S_{11} & D_{12} & D_{15} \\ S_{11} & D_{13} & D_{15} \end{pmatrix} + \Delta$$

 (N, Δ)

H. O.

$$\omega=2 \quad \text{---} 2^+_S \text{---} 2^+_M \text{---} 1^+_A \text{---} 0^+_M \text{---} 0^+_S$$

$$\omega=1 \quad \text{---} 1^-_M$$

$$\omega=0 \quad \text{---} 0^+_S$$



$\frac{E_0 + 2\Omega}{1_A} 1_A^+$

$-\frac{\Delta}{3} 2_M^+$

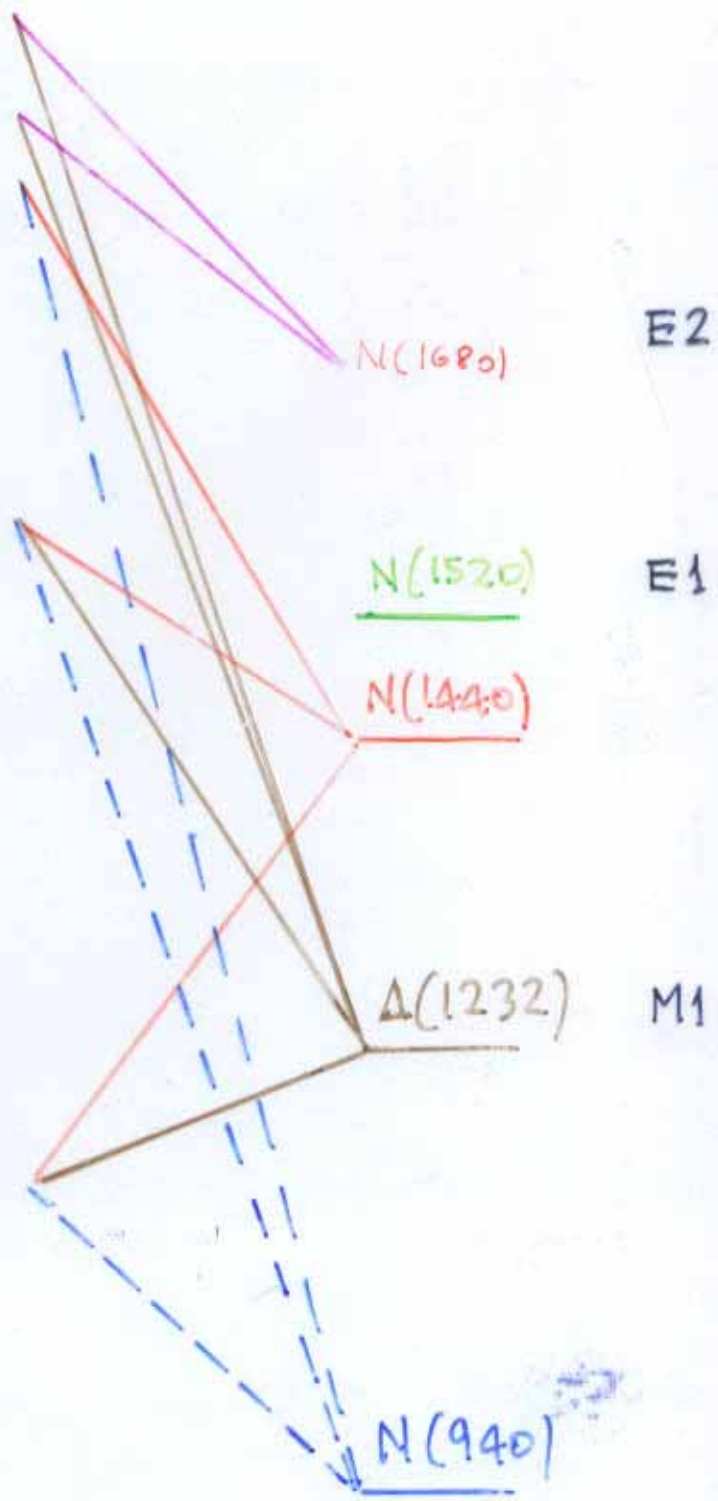
$-\frac{2\Delta}{3} 2_S^+$

$-\frac{\Delta}{2} 0_M^+$

$\frac{2\hbar\omega}{(0_S^+, 0_M^+, 2_S^+, 2_M^+, 1_A^+)$

$\frac{1\hbar\omega}{(1_M^-)} -\Delta 0_S^+, 1_M^-$

$\frac{0\hbar\omega}{0_S^+} E_0 0_S^+$



$T + V_{h.o.} + U + H_{hyp} \leftarrow OGE$

two analytical solutions

A) harmonic oscillator

$$\sum_{i < j} \frac{1}{2} K (\vec{r}_i - \vec{r}_j)^2 = \frac{3}{2} K x^2$$

$$E = (3 + N) \hbar \omega$$

$$\hookrightarrow 2\nu + \gamma$$

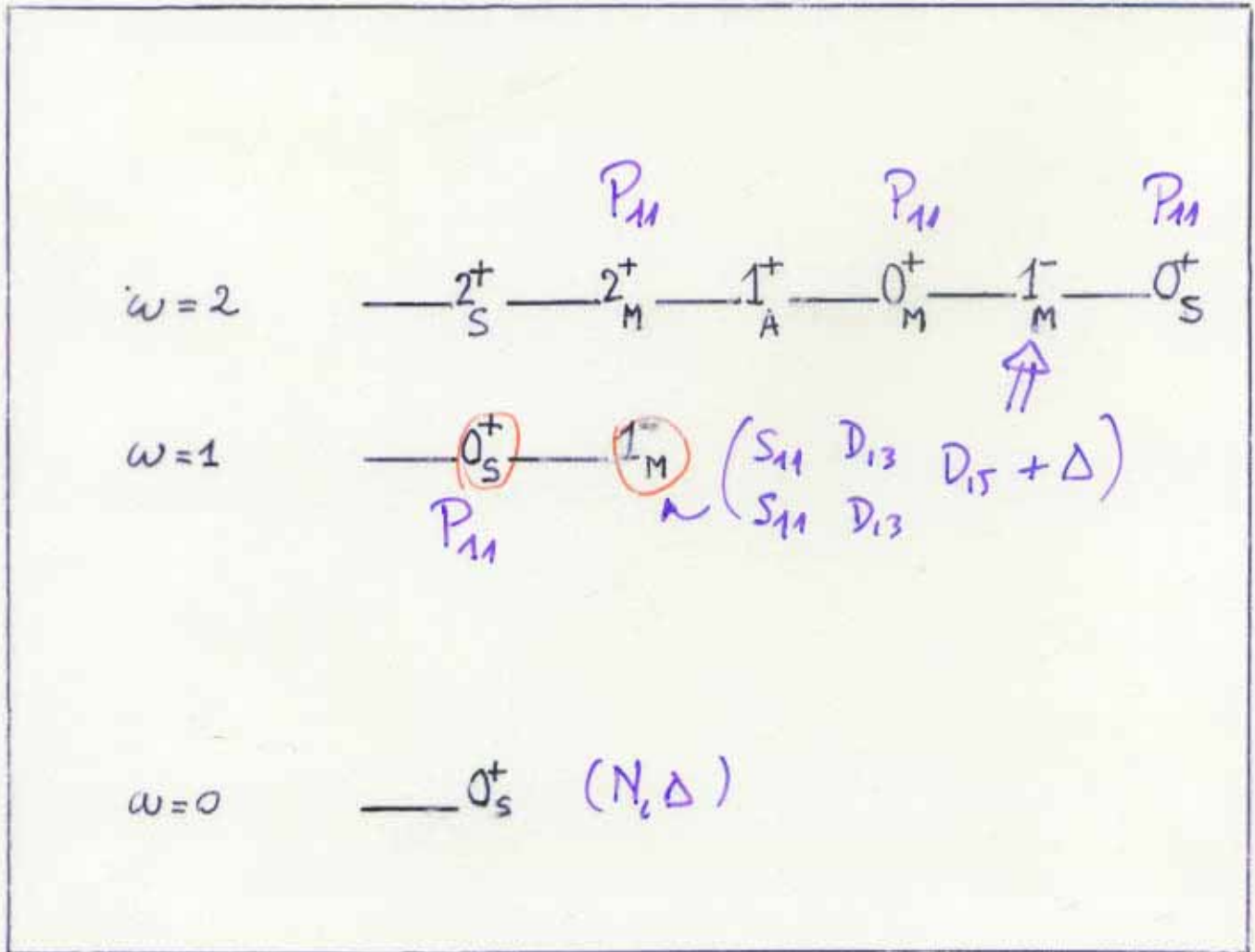
B) hypercoulomb

$$V_{\text{hyc}}(x) = -\frac{\tau}{x}$$

$$E = -\frac{m \tau^2}{2 N^2}$$

$$\hookrightarrow \frac{5}{2} + \gamma + \nu$$

$-\frac{2}{x}$ "hyperCoulomb"



(M.M. Giannini, F. Iachello, E. Santopinto)

• symmetry

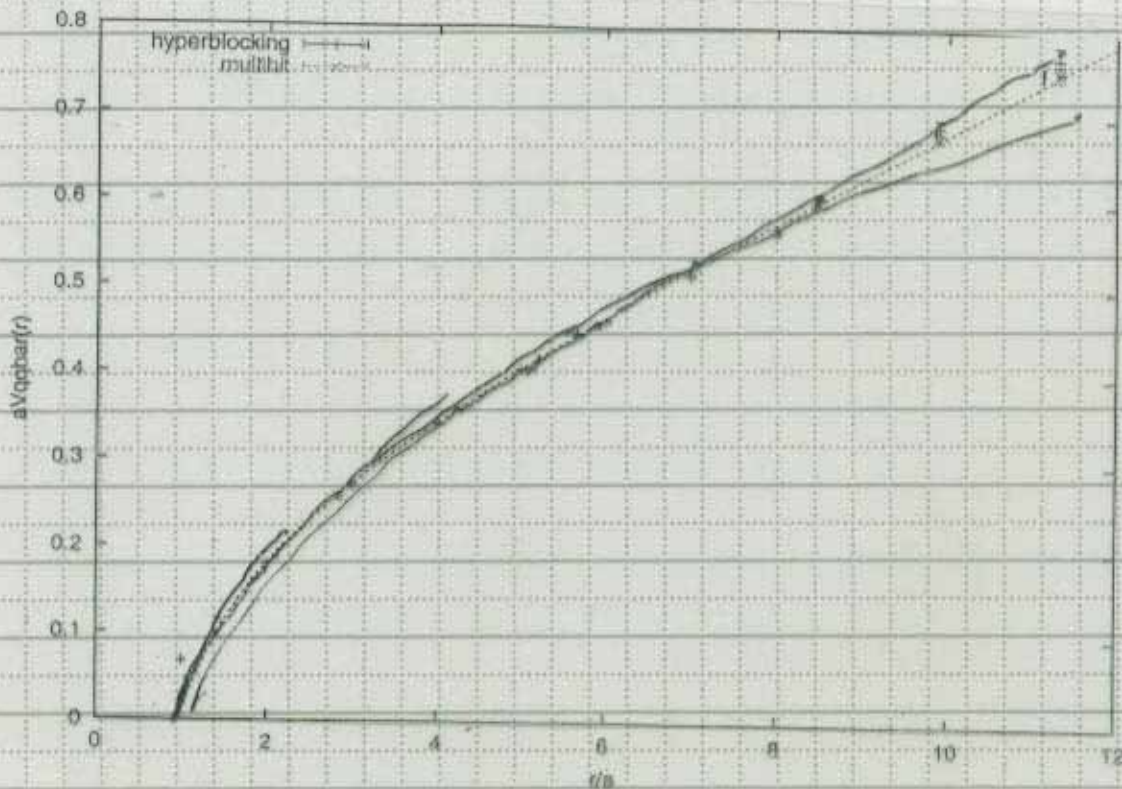
S_3 :

S (symmetric)

M (mixed symmetric)

A (antisymmetric)

LATTICE QCD $q\bar{q}$ potential



$$V_{q\bar{q}}(r) = -\frac{A_{q\bar{q}}}{r} + \sigma_{q\bar{q}} r + C_{q\bar{q}}$$

(Cornell pot.)

G. Bali et al.

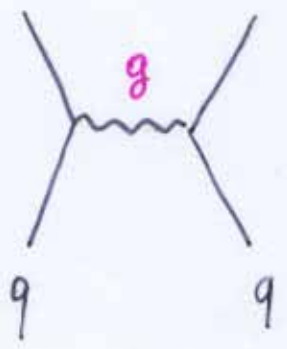
Phys. Rev. D51, 5165 (1995)

HYPERCENTRAL MODEL

$$V_{3q} = V(x) + H_{hyp}$$

- $V(x) = -\frac{\tau}{x} + \alpha x$ $x = \sqrt{\rho^2 + \lambda^2}$

- $H_{hyp} = \sum_{i < j} V^s(\vec{r}_i, \vec{r}_j) \vec{\sigma}_i \cdot \vec{\sigma}_j + \text{tensor}$
 $\quad \quad \quad \downarrow$
 $\quad \quad \quad \delta(\vec{r}_i - \vec{r}_j)$



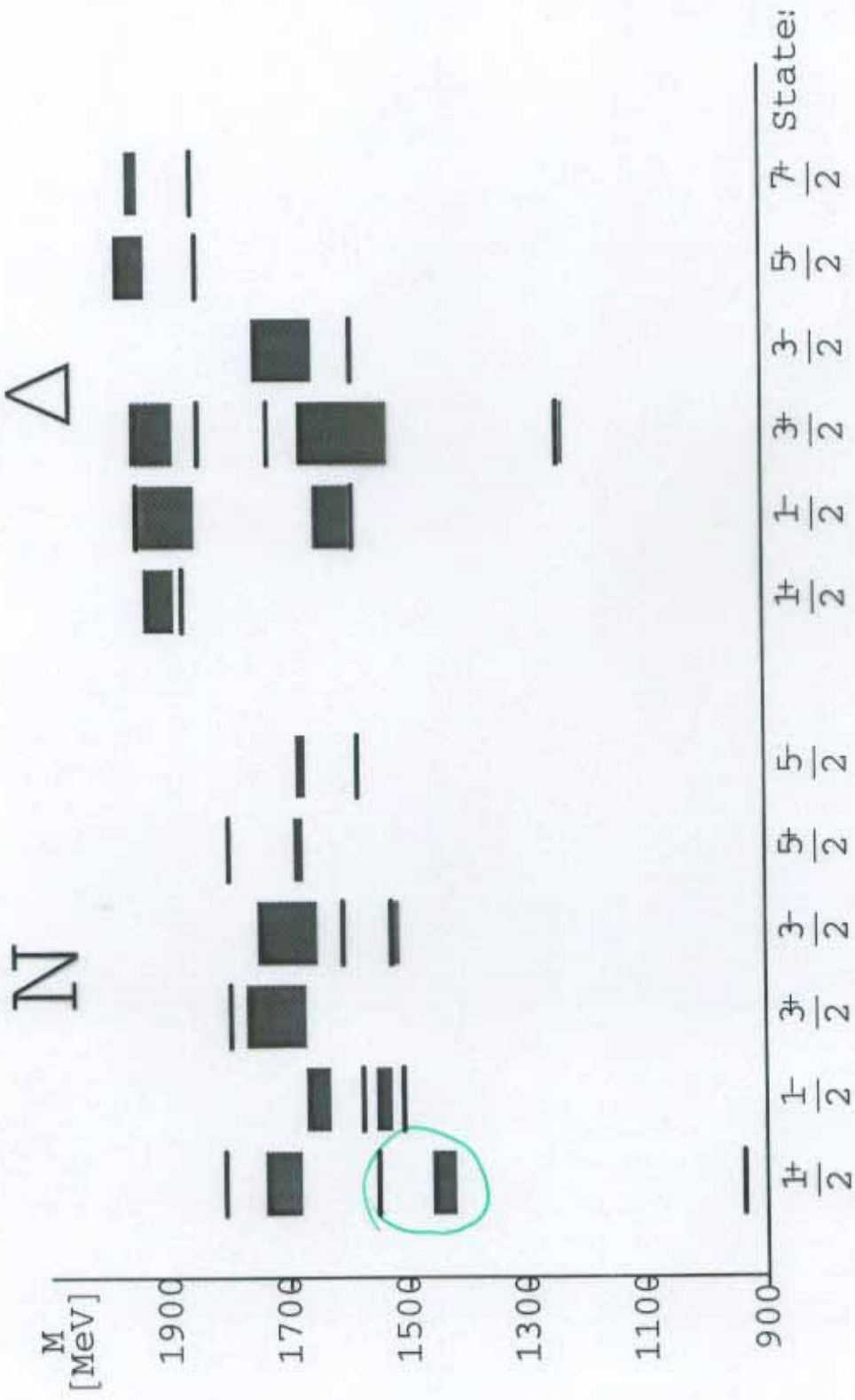
3 parameters

$$\rightarrow \langle r_p^2 \rangle^{1/2} \sim 0.48 \text{ fm}$$

τ
 α
 strength of H_{hyp}

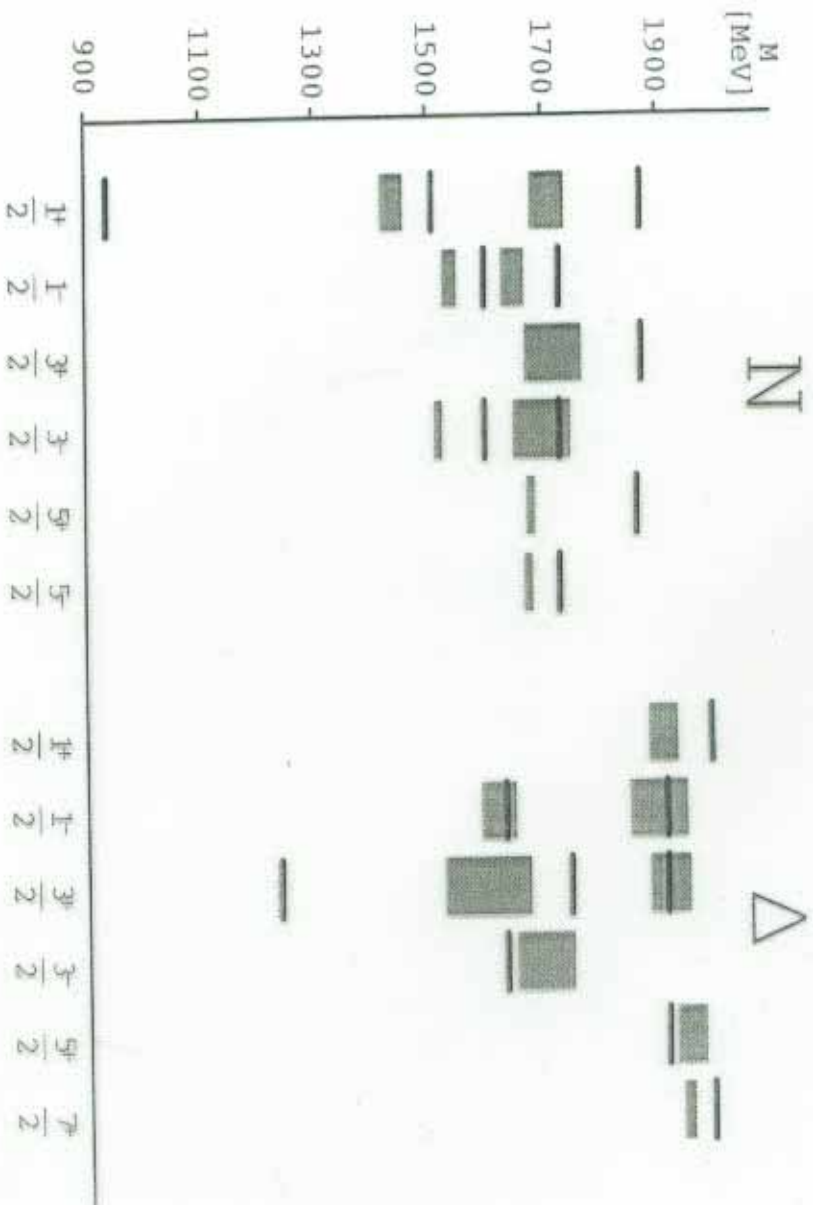
← fixed to the spectrum

$$\left[m = \frac{M}{3} \right]$$



M. Ferraris et al., Phys. Lett. B 364 (1995) 231

h CQM with $T = \sum_{i=1}^3 \sqrt{m^2 + p_i^2}$
 (H_{free})



Isospin dependence

$$\mathcal{H} = \mathcal{H}_0 + \mathcal{H}_S + \mathcal{H}_I + \mathcal{H}_{SI} + V_T$$

- $\mathcal{H}_0 = T - \frac{\tau}{x} + \alpha x$

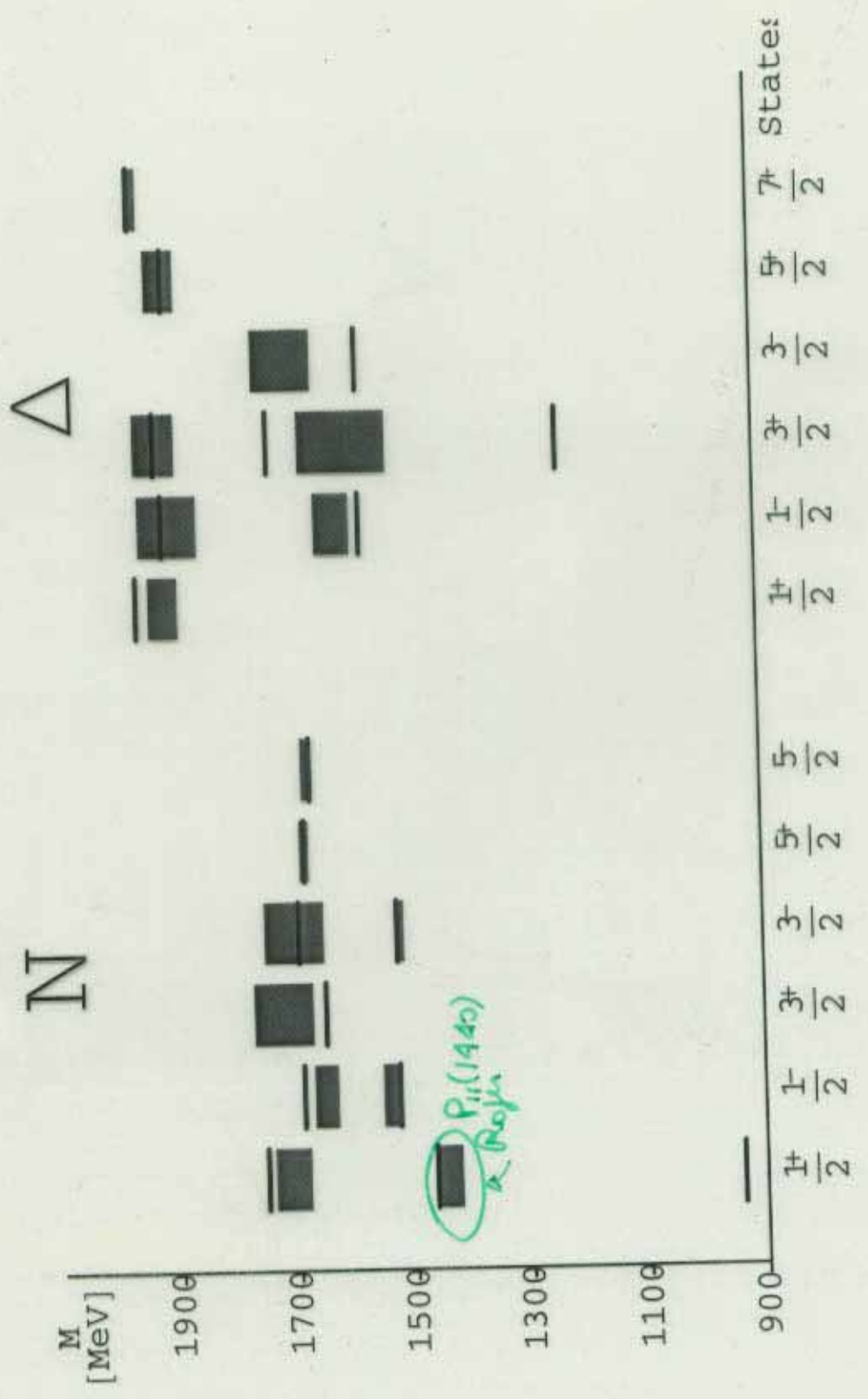
- $\mathcal{H}_S = A_S \sum_{i < j} f_S(\mathbf{r}_{ij})(\mathbf{S}_i \cdot \mathbf{S}_j)$

- $\mathcal{H}_I = A_I \sum_{i < j} f_I(\mathbf{r}_{ij})(\mathbf{I}_i \cdot \mathbf{I}_j)$

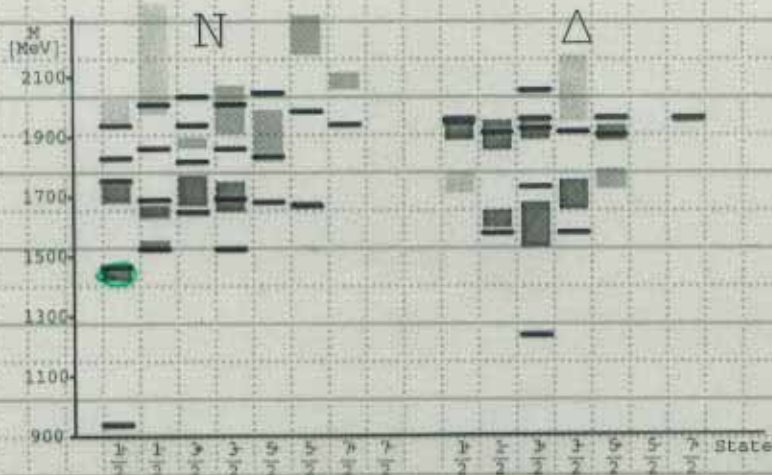
- $\mathcal{H}_{SI} = A_{SI} \sum_{i < j} f_{SI}(\mathbf{r}_{ij})(\mathbf{S}_i \cdot \mathbf{S}_j)(\mathbf{I}_i \cdot \mathbf{I}_j)$

$$- f_k(\mathbf{r}_{ij}) = \frac{1}{(\sqrt{\pi}\sigma_k)^3} e^{-\frac{r_{ij}^2}{\sigma_k^2}} \quad k = S, I, SI$$

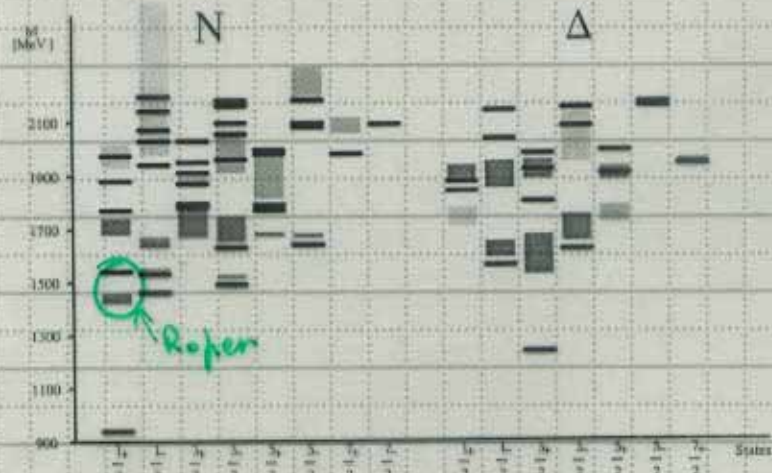
- $V_T = \sum_{i < j} \frac{2\alpha_s}{3m^2} \left\{ \frac{1}{r_{ij}^3} \left[3 \frac{(\mathbf{S}_i \cdot \mathbf{r}_{ij})(\mathbf{S}_j \cdot \mathbf{r}_{ij})}{r_{ij}^2} - \mathbf{S}_i \cdot \mathbf{S}_j \right] \right\}$



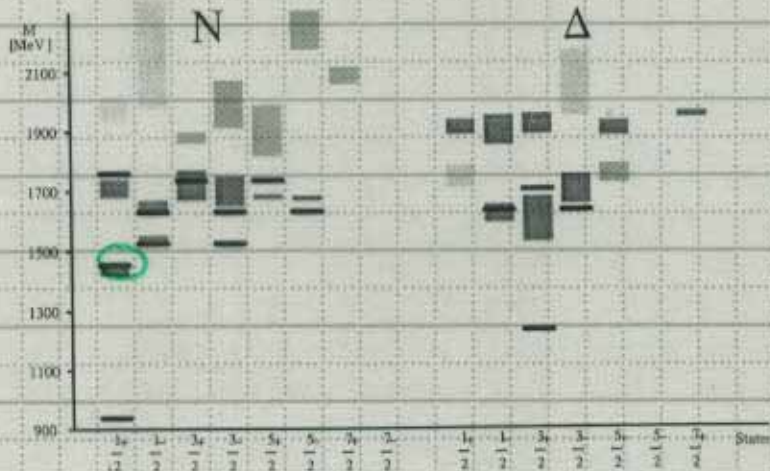
M.G., E. Santopinto, A. Vassallo, EPJ A12, 447 (2001)



HCQM



Capstick-Isgur



Plassas et al.

$M(\Delta) - M(N) :$

$H_{hyp} \rightarrow 35\%$

$H_I \rightarrow 15\%$

$H_{SI} \rightarrow 50\%$

PDG

h CQM

h CQM + 150

 S_{11}

1535

1507

1524

1650

1574

1688

1887

1861

(2090)

1937

2008

 D_{13}

1520

1526

1524

1700

1606

1692

1899

1860

(2080)

1969

2008

 P_{13}

1720

1797

1628

1900

1835

1816

1853

1894

1867

1939

2034

 P_{33}

1232

1240

1232

1600

1727

1727

1823

1921

1920

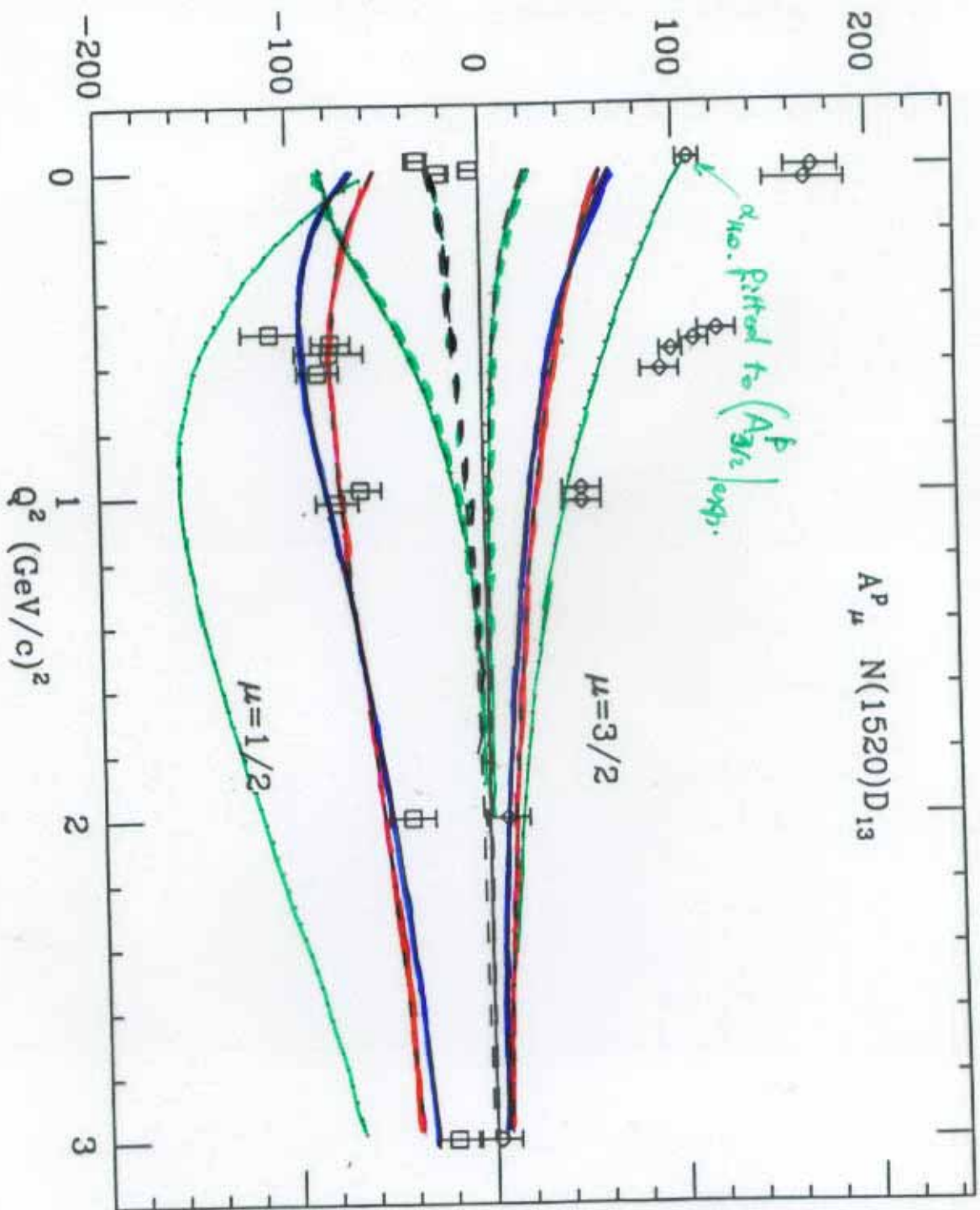
1856

1953

2104

2049

$10^{-3} \text{ GeV}^{-1/2}$



$A_{\mu}^p \ N(1520)D_{13}$

— h.e.
— SIG= HCQM analysis.
— HCQM (Gervais)

Aiello et al. J. Phys. G124, 753 (1998)

Helicity Amplitudes (Photocouplings)

$$\sigma + N \rightarrow N^*, \Delta$$

$$A_{\frac{1}{2}} = \langle B, J', J'_z = \frac{1}{2} | H_{em}^{\epsilon} | N, J = \frac{1}{2}, J_z = -\frac{1}{2} \rangle$$

$$A_{\frac{3}{2}} = \langle B, J', J'_z = \frac{3}{2} | H_{em}^{\epsilon} | N, J = \frac{1}{2}, J_z = +\frac{1}{2} \rangle$$

$$H_{em}^{\epsilon} = - \sum_{i=1}^3 \left[\frac{e_i}{2m_i} (\vec{p}_i \cdot \vec{A}_i + \vec{A}_i \cdot \vec{p}_i) + 2\mu_i \vec{S}_i \cdot (\vec{\nabla} \wedge \vec{A}_i) \right]$$

e_i quark charge

m_i " mass

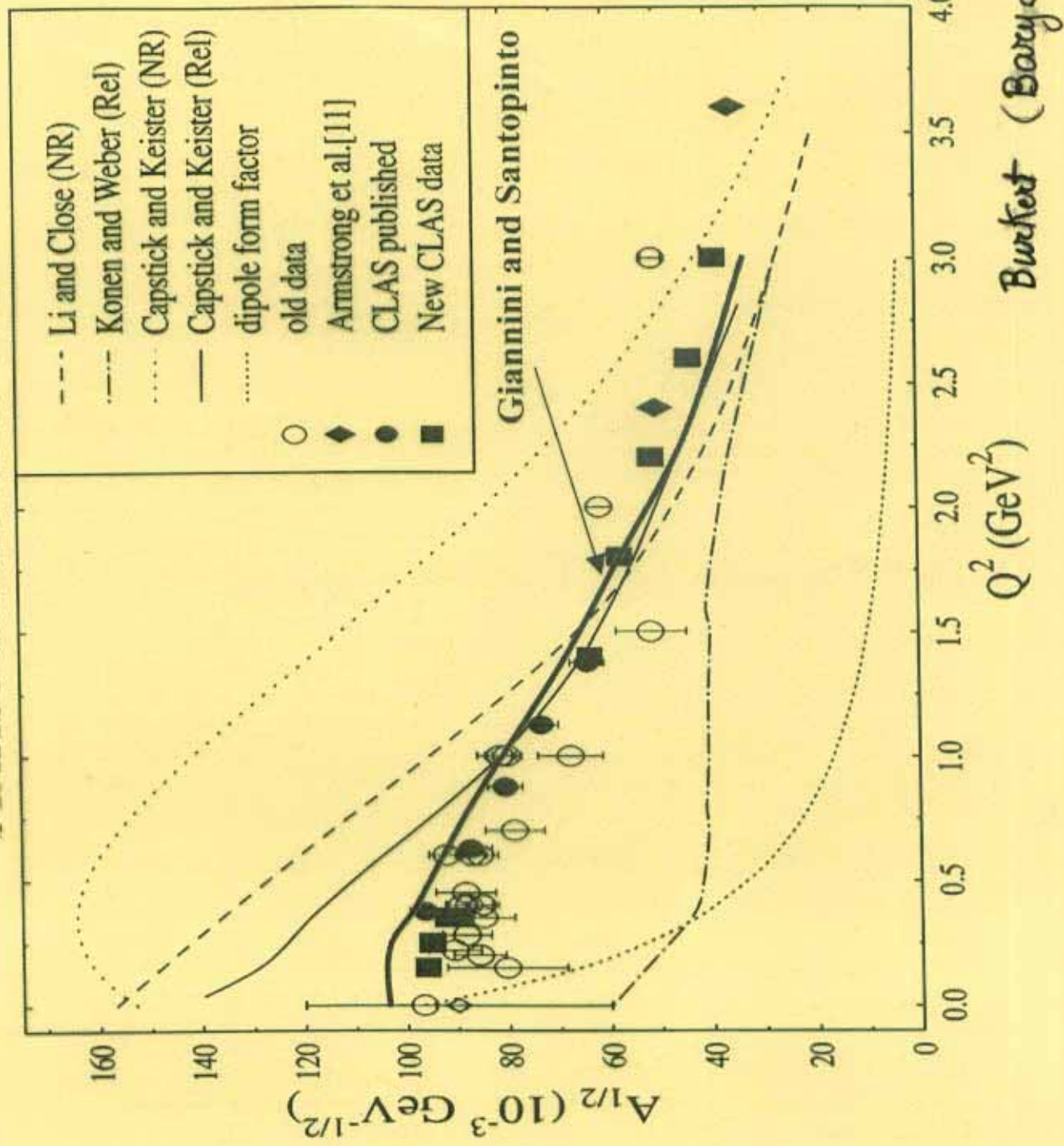
μ_i " magnetic moment

S_i " spin

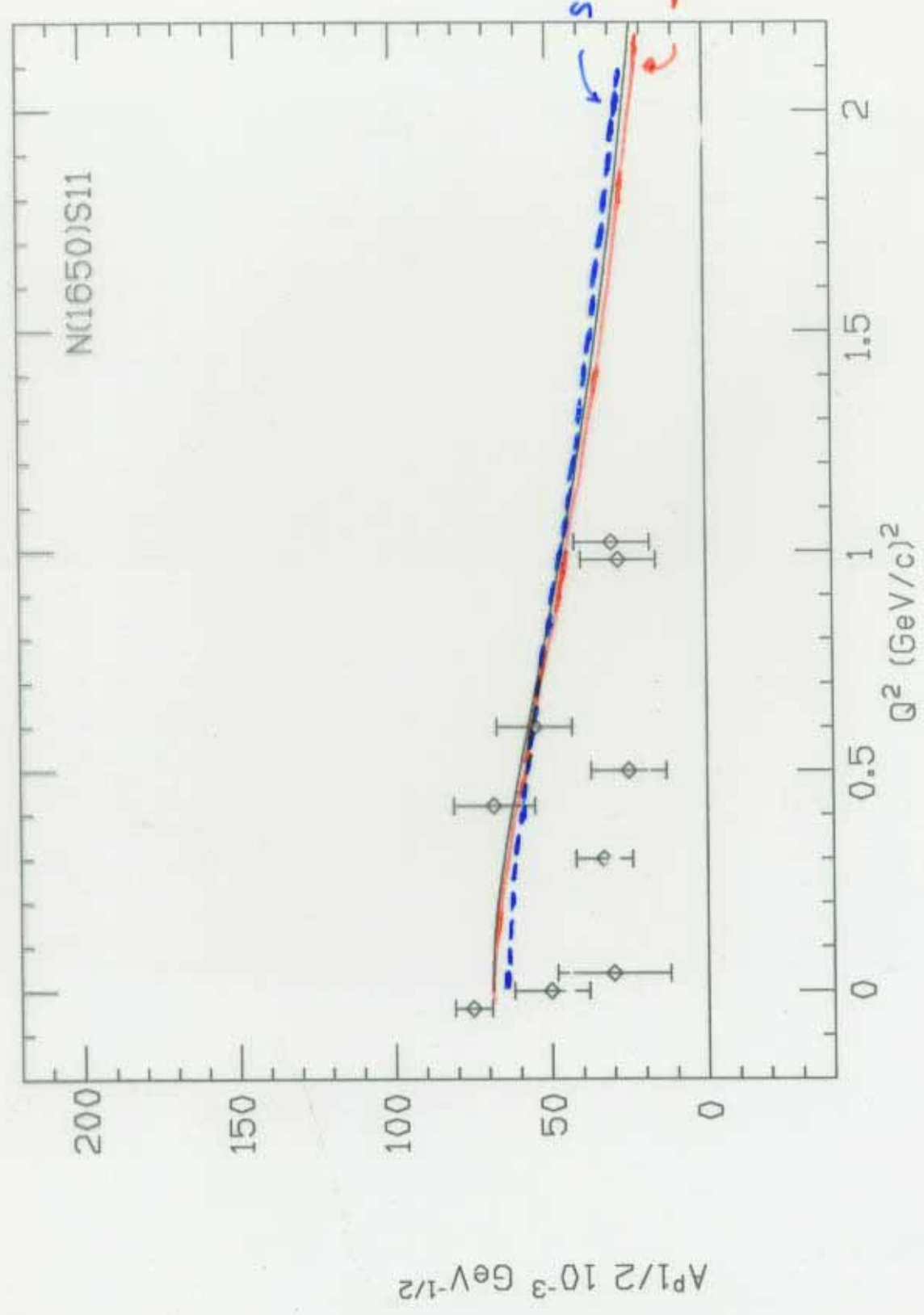
p_i " momentum

$\vec{A}_i = \vec{A}(\vec{r}_i)$ photon field

Photocoupling amplitude $A_{1/2}$

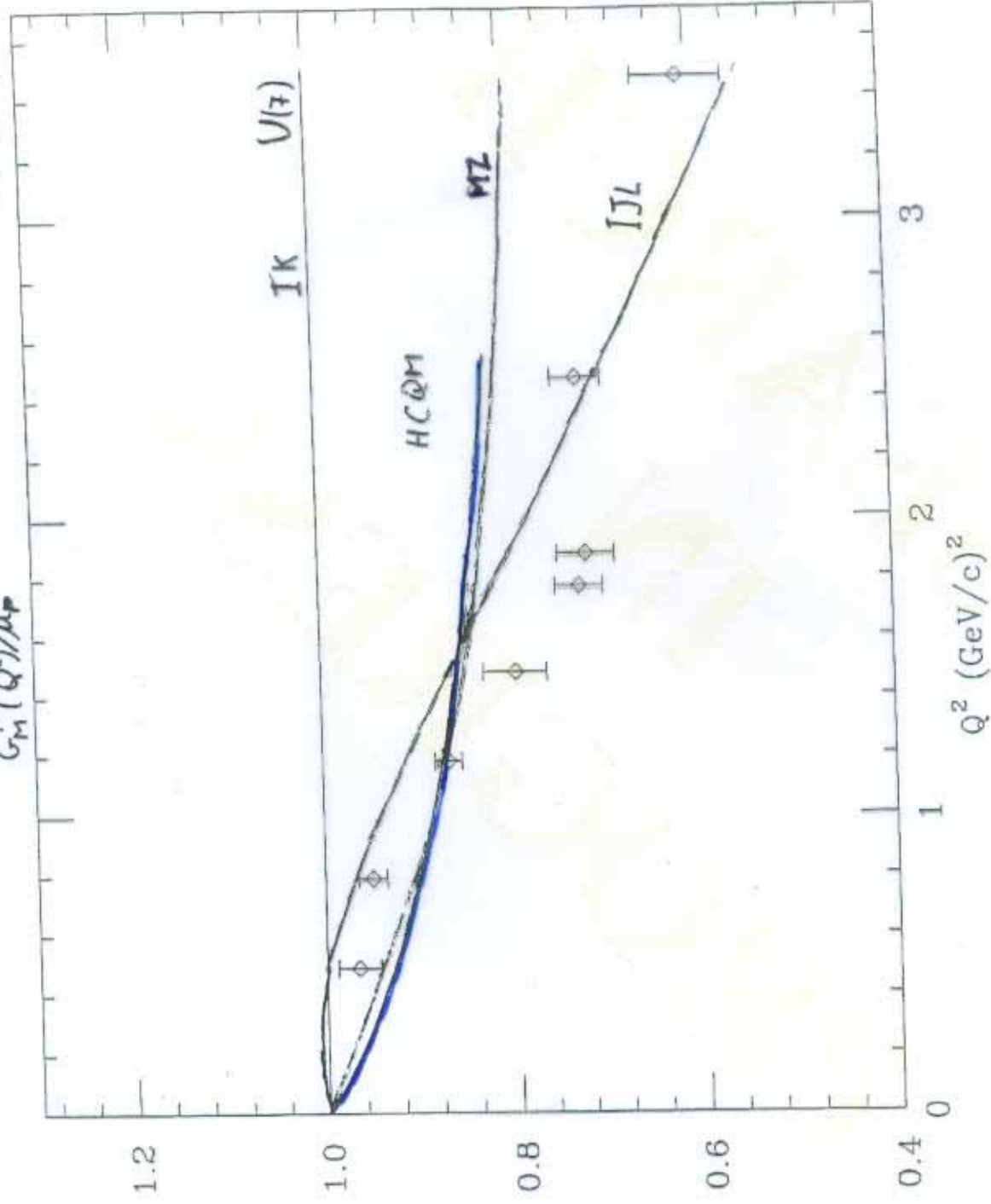


J. Phys. G: 24,
753 (1998)



$$R = \frac{G_E^p(Q^2)}{G_M^p(Q^2)/\mu_p}$$

PRC 62, 025208 (2000)



T_{NR} , boosts, current expansion

Geneva, Sept. 2002

