

Helicity dependence of the $N\pi(\pi)$ photoproduction: new results from Mainz

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for the GDH and A2 collaborations



Nstar 2002 Workshop - Pittsburgh

SUMMARY

- Physical motivations
- Experimental setup
- Results [$\vec{\gamma} \vec{p} \rightarrow N\pi(\pi)$]
- Future outlook [$\vec{\gamma} \vec{n} \rightarrow N\pi(\pi)$]
(with some results)
- Conclusions

The GDH collaboration

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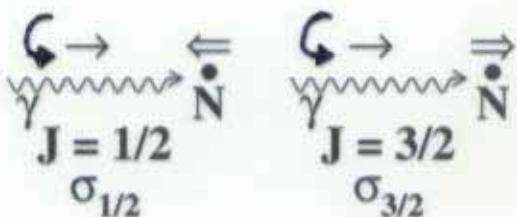
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Experimental Programme

- Experimental check of the **Gerasimov-Drell-Hearn** sum rule both for **proton** and **neutron** at
 - Mainz (MAMI) ($m_\pi \leq E_\gamma \leq 800$ MeV)
 - Bonn (ELSA) (600 MeV $\leq E_\gamma \leq 3$ GeV)
- Measurement at **Mainz** of the **helicity dependence** of all the
 - $\gamma N \rightarrow N\pi$
 - $\gamma N \rightarrow N\pi\pi$channels as a powerful tool to study the properties of baryon resonances ($\Delta, P_{11}, D_{13}, S_{11}$)

Gerasimov-Drell-Hearn sum rule

- derived by S.B.Gerasimov (1966) – S.D.Drell, A.C.Hearn (1966)
- absorption of circularly polarized photons by longitudinally polarized nucleons



- causality (unsubtracted dispersion relation), optical theorem, low-energy theorems

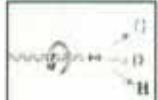
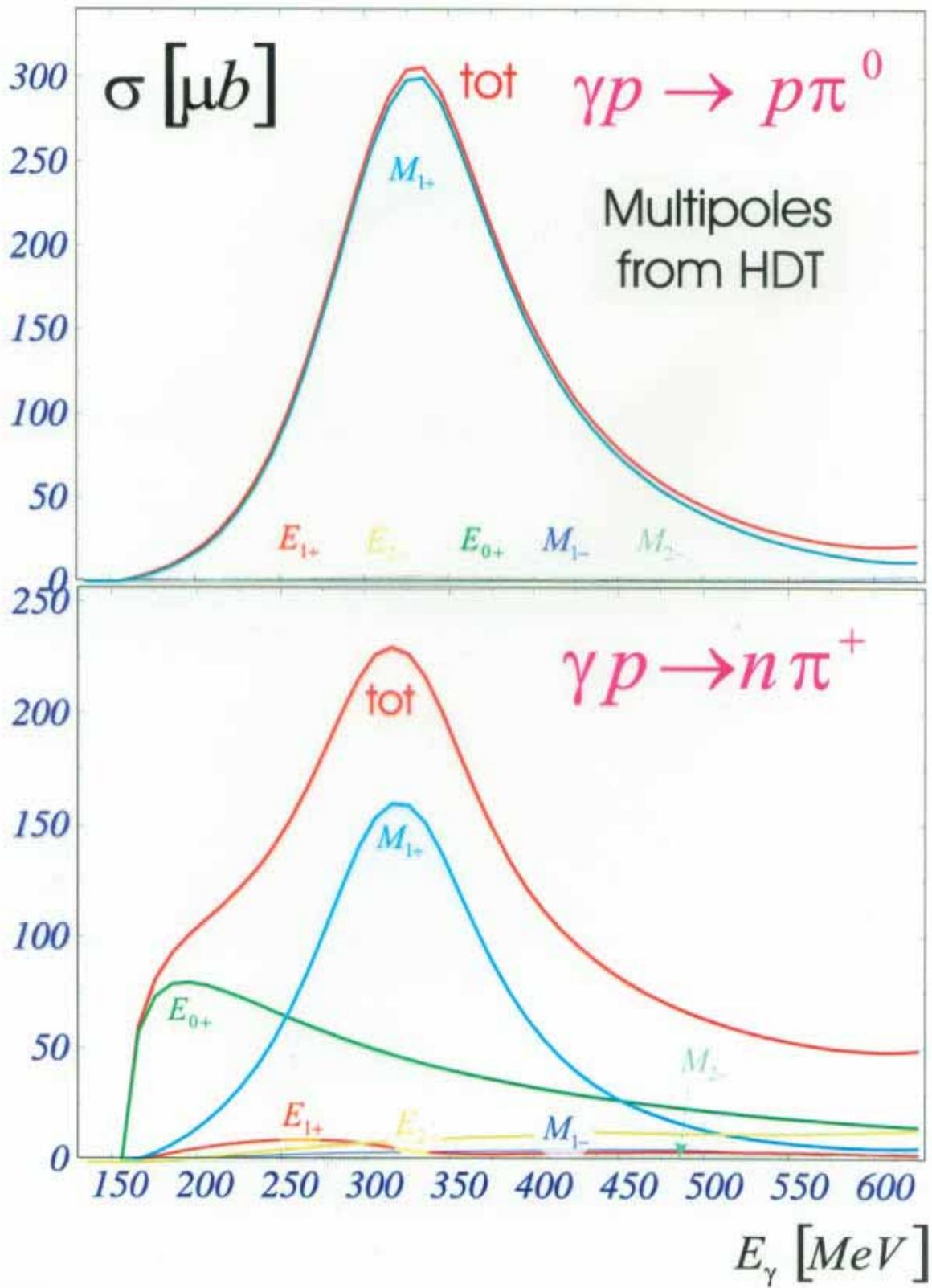
$$\int_{m_\pi}^{\infty} \frac{(\sigma_{3/2} - \sigma_{1/2})}{E_\gamma} dE_\gamma = \frac{2\pi^2 \alpha}{m^2} \kappa^2$$

- fundamental check of our knowledge of the γN interaction
- important comparison test for nucleon models
- “spin polarizability” γ_0 :

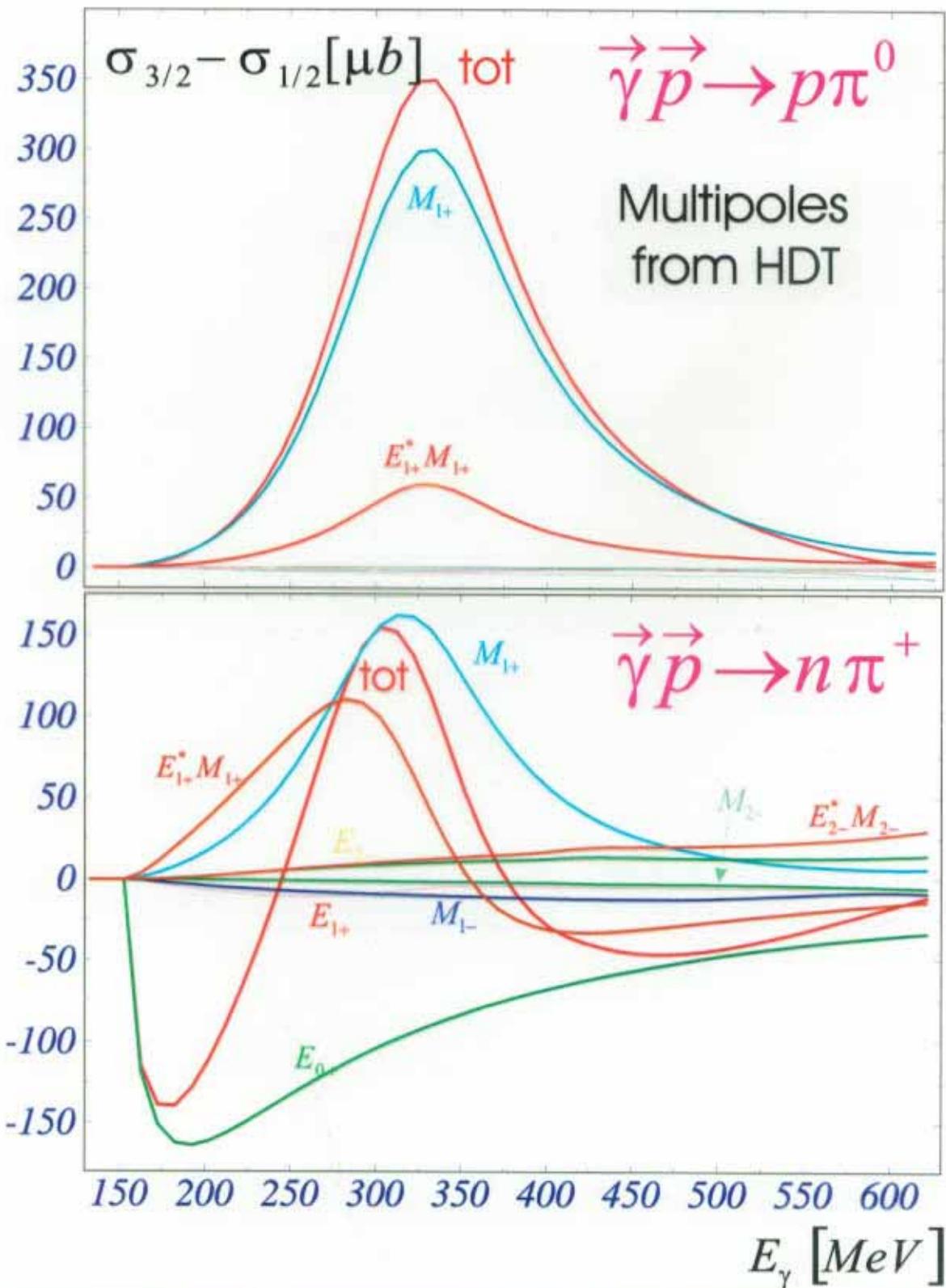
$$\gamma_0 = -\frac{1}{4\pi} \int_{m_\pi}^{\infty} \frac{(\sigma_{3/2} - \sigma_{1/2})}{E_\gamma^3} dE_\gamma$$



Motivation



Motivation



Helicity amplitudes

$$\begin{aligned} \mathbf{A}_{1/2} &= \frac{1}{\sqrt{2q}} \langle N^*(J, M = 1/2) | \mathbf{J}_+ | N(J_i = 1/2, M_i = -1/2) \rangle \\ \mathbf{A}_{3/2} &= \frac{1}{\sqrt{2q}} \langle N^*(J, M = 3/2) | \mathbf{J}_+ | N(J_i = 1/2, M_i = +1/2) \rangle \end{aligned}$$

q : photon momentum

\mathbf{J}_+ : hadronic current corresponding to the absorption of a γ with positive helicity on a nucleon N with spin J_i and spin projection M_i leading to a resonance state N^* with spin $J \geq 1/2$ and spin projection M

- \mathbf{J}_+ strongly depends in the quark-quark interactions (\Rightarrow NPQCD)
- helicity amplitudes related to e.m. multipoles, e.g.

$$\begin{aligned} \mathbf{M}_{1+} &= -\frac{1}{2\sqrt{3}} (\sqrt{3}\mathbf{A}_{1/2} + 3\mathbf{A}_{3/2}) \\ \mathbf{E}_{1+} &= -\frac{1}{2\sqrt{3}} (\sqrt{3}\mathbf{A}_{1/2} - \mathbf{A}_{3/2}) \end{aligned}$$

- $\sigma_{unpol} \propto |\mathbf{A}_{1/2}|^2 + |\mathbf{A}_{3/2}|^2$

Helicity Amplitudes (Proton target)

(units of $10^{-3} \text{ GeV}^{-1/2}$)

	P_{33} (1232)		P_{11} (1440)	D_{13} (1520)		S_{11} (1535)
	$A_{1/2}^p$	$A_{3/2}^p$	$A_{1/2}^p$	$A_{1/2}^p$	$A_{3/2}^p$	$A_{1/2}^p$
NRCQM (Isgur-Karl)	-103	-179	-24	-23	+128	+147
NRCQM (Sartor-Stancu)	-101	-181	-31	+45	+202	+203
RCQM (Close-Li)	-113	-195	+10	-30	+146	+163
NRCQM+Mix. (Close-Li)	-94	-162	-93	-51	+133	+142
NRCQM+RC (Pfeil <i>et al.</i>)	-93	-163	-105	-79	+65	+116
RCQM (Pfeil <i>et al.</i>)	-81	-170	-9	-7	+63	+54
RCQM (Kapstick)	-108	-186	+4	-15	+134	+76
Light Front (Kapstick)	-100	-180	+35	+25	+200	+90
Algebraic appr. (Iachello <i>et al.</i>)	-91	-157	0	-43	+109	+162
NRCQM+Mix. (Giannini <i>et al.</i>)	-93	-162	-69	-47	+85	+116
P.W. Analyses	[-147,-128]	[-264,-247]	[-129,-58]	[-43,-5]	[75,206]	[29,125]
PDB best value	-135 ± 6	-255 ± 8	-65 ± 4	-24 ± 9	$+166 \pm 5$	$+90 \pm 30$

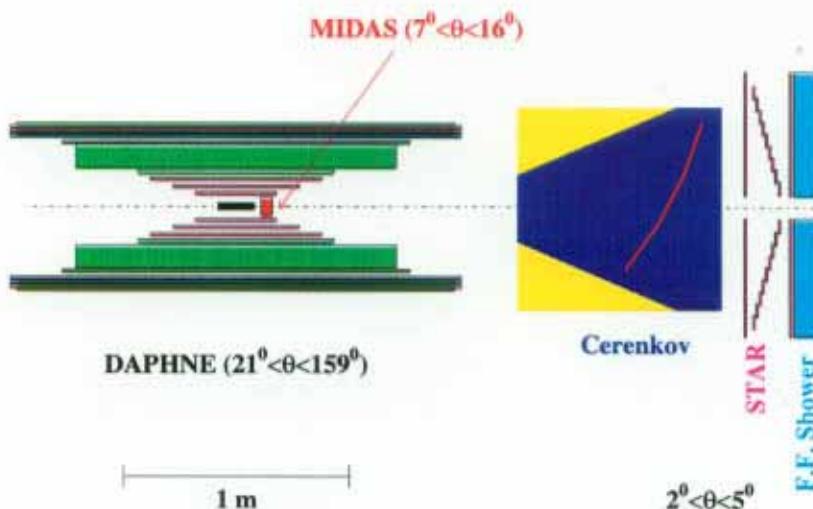
PREVIOUS
VALUE

Setup

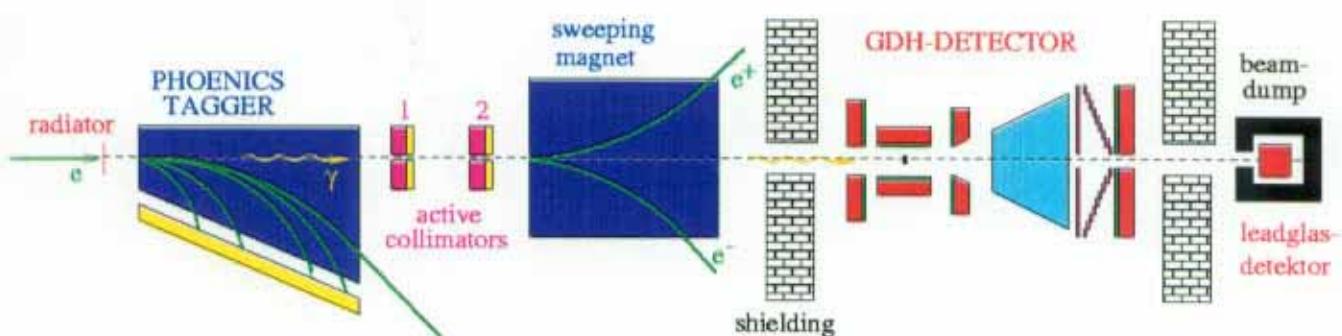
Experimental set-up

- Tagged photon facility of the MAMI/ELSA accelerators
- circularly polarized photons \Rightarrow helicity transfer from bremsstrahlung of linearly polarized electrons (strained GaAs source)
- Longitudinally polarized nucleons
 - “frozen spin” (deuterated) butanol $[C_4 H(D)_9 OH(D)]$
 - ^6LiD target (future: ^3He gas target)

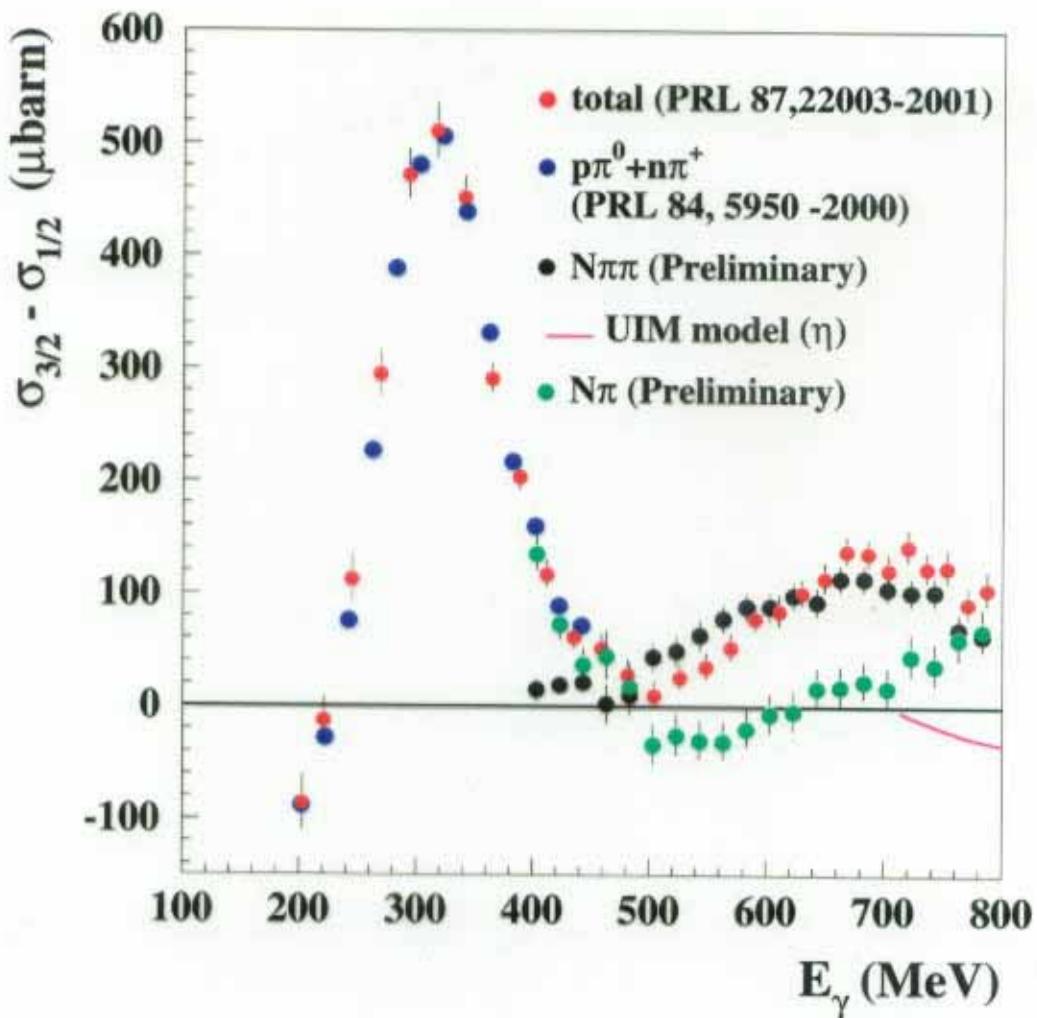
The GDH Detector at MAMI



The GDH detector at Bonn

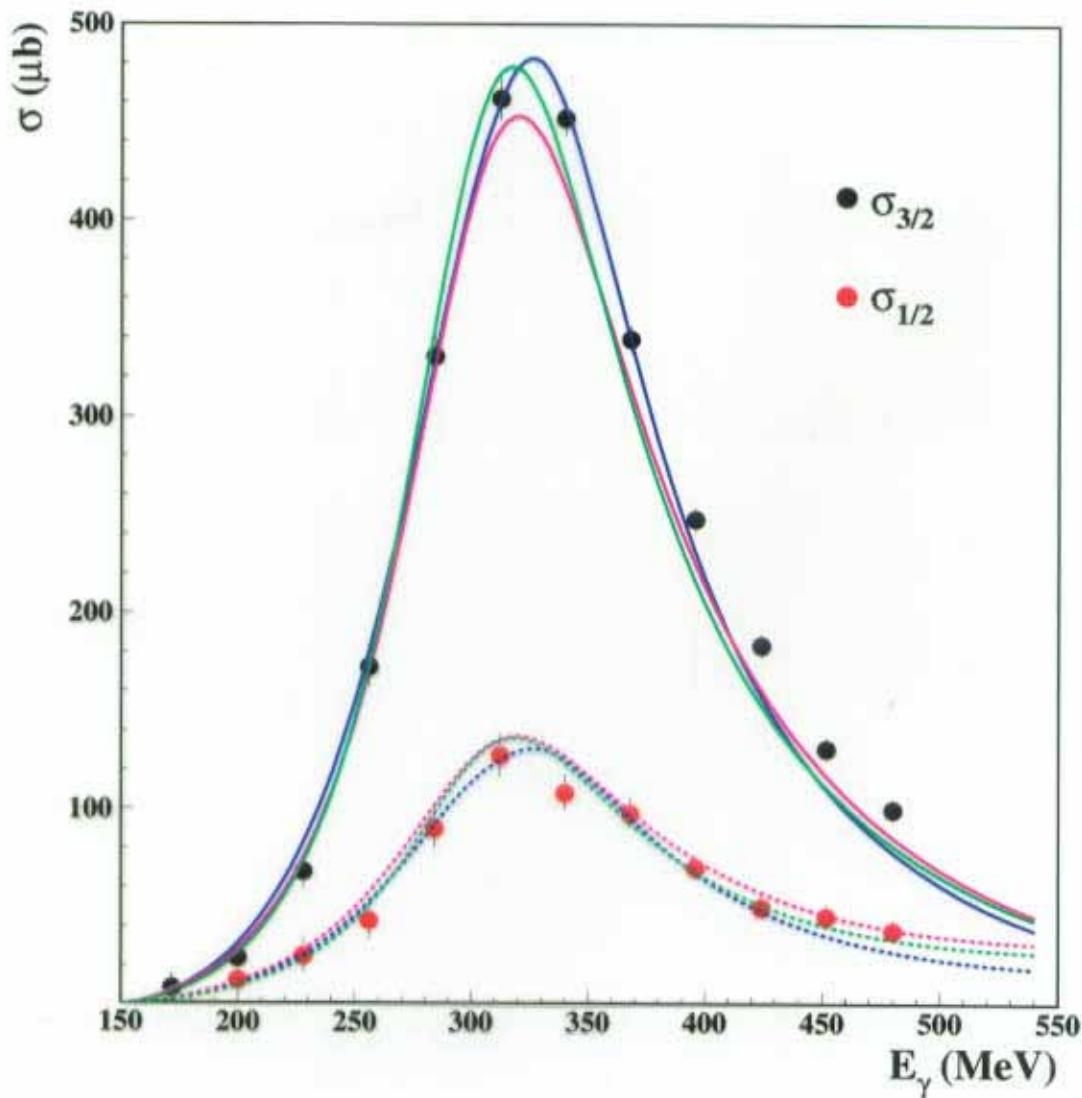


Results

 $\vec{\gamma} \vec{p} \rightarrow \text{hadrons}$ 

- ($N\pi$) cross section above 450 MeV = (total) - ($N\pi\pi$) - ($N\rho$)

$$\vec{\gamma} \vec{p} \rightarrow p\pi^0$$

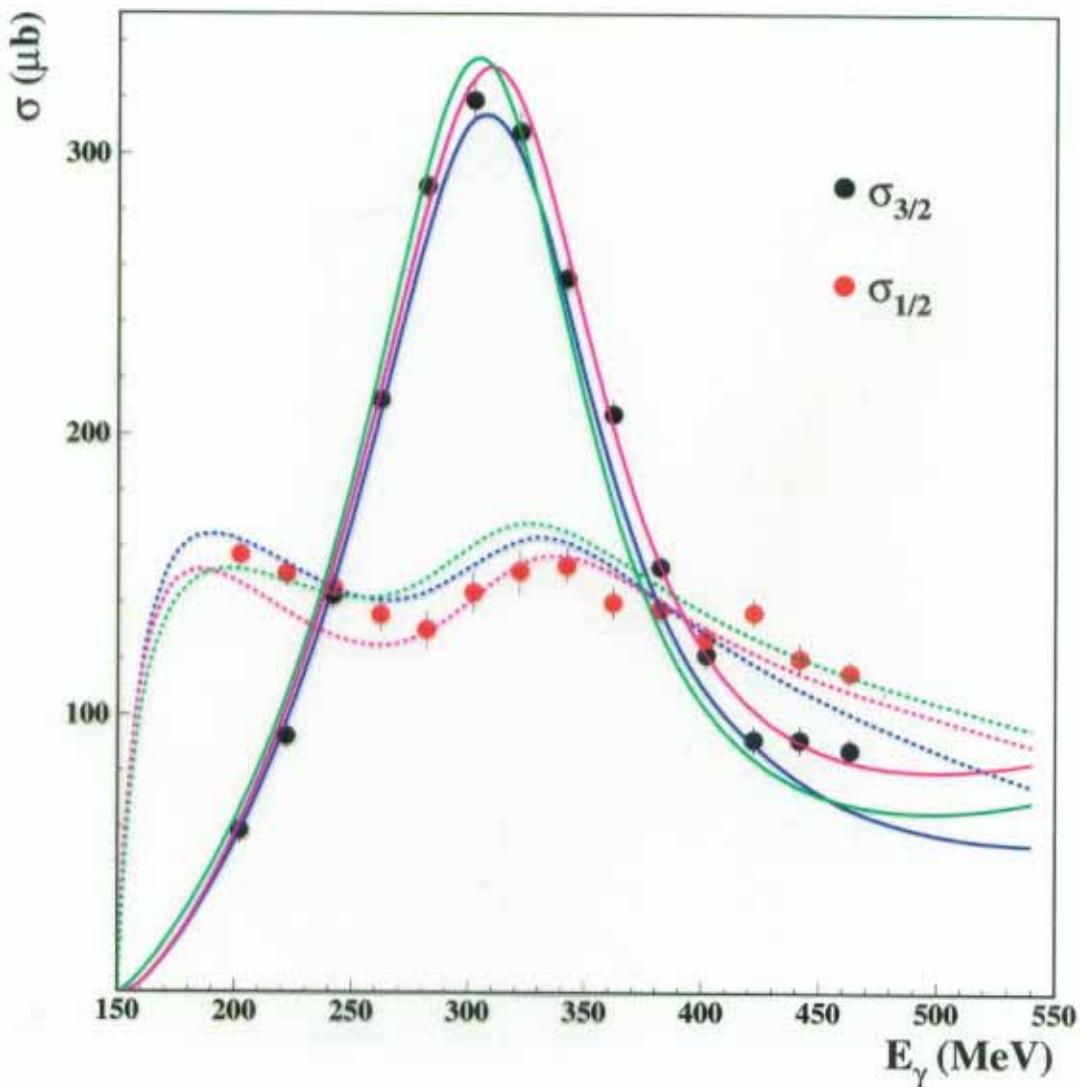


HDT: dispersion theory, Hanstein et al., NPA 632, 521 (99)

SAID: multipole analysis (solution SM02K), R.Arndt et al., nucl-th/0205067

UIM (MAID): multipole analysis, Drechsel et al., PRD 59, 094021 (99)

$$\vec{\gamma} \vec{p} \rightarrow n\pi^+$$

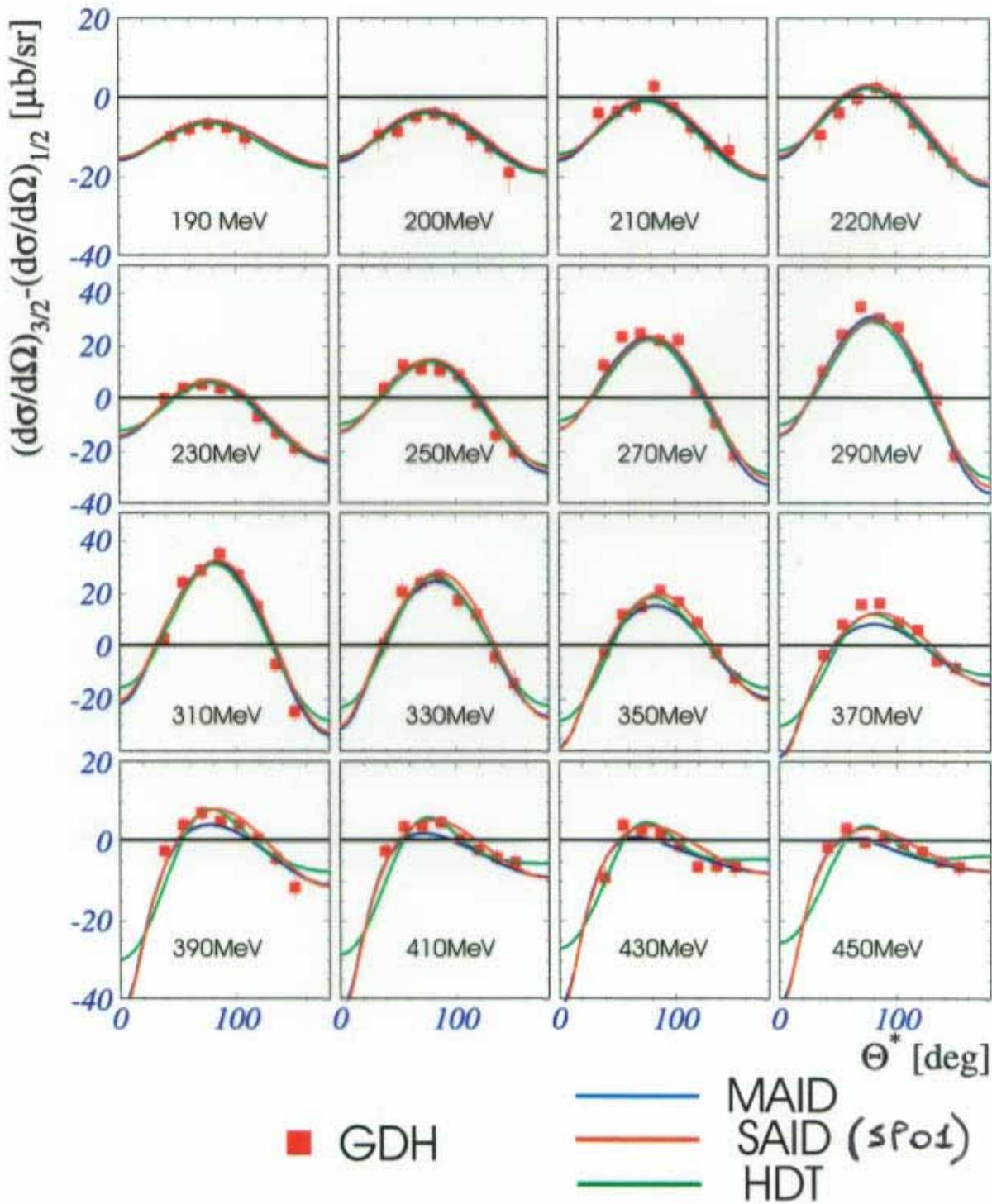


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Results

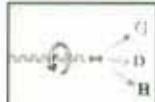
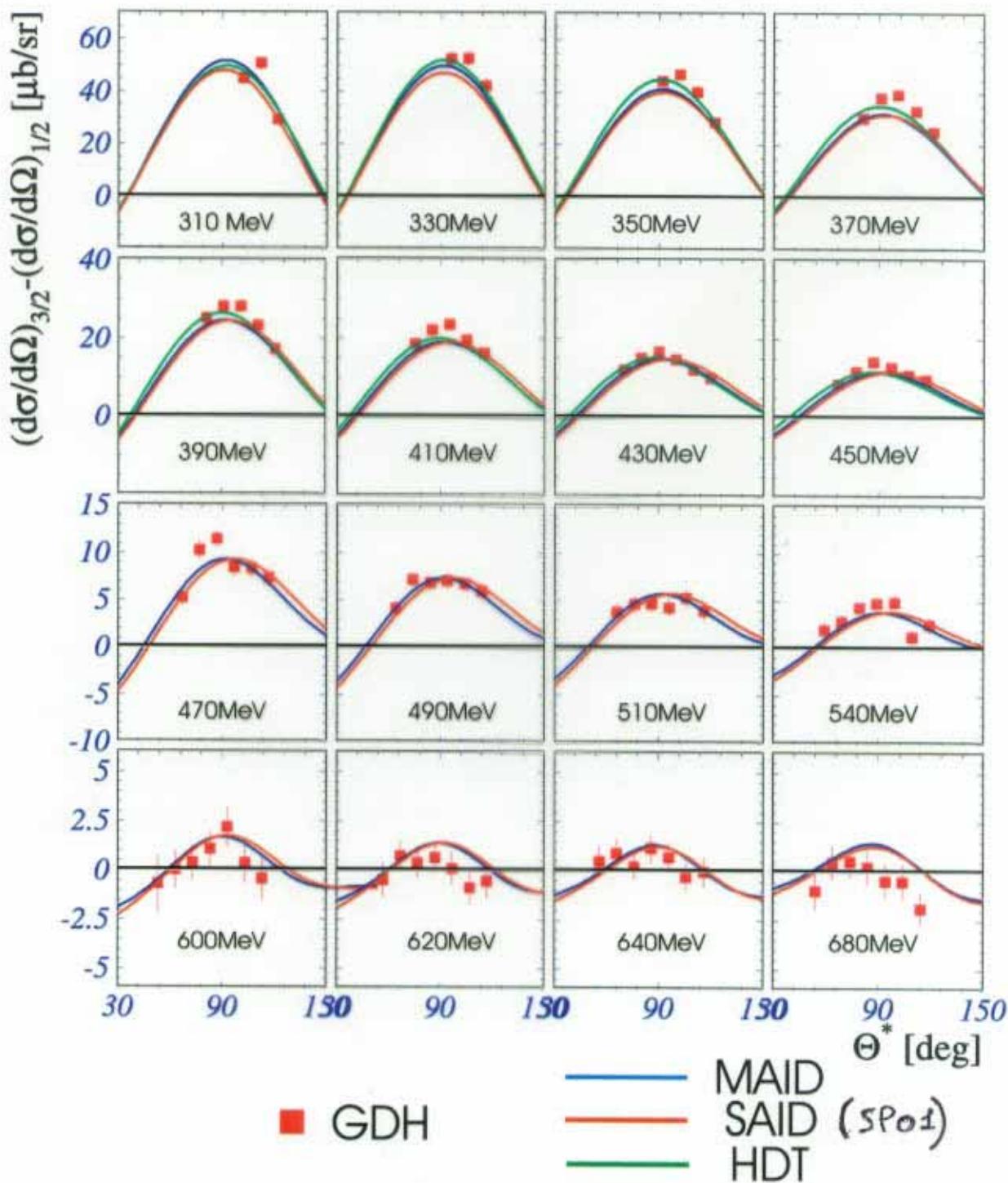


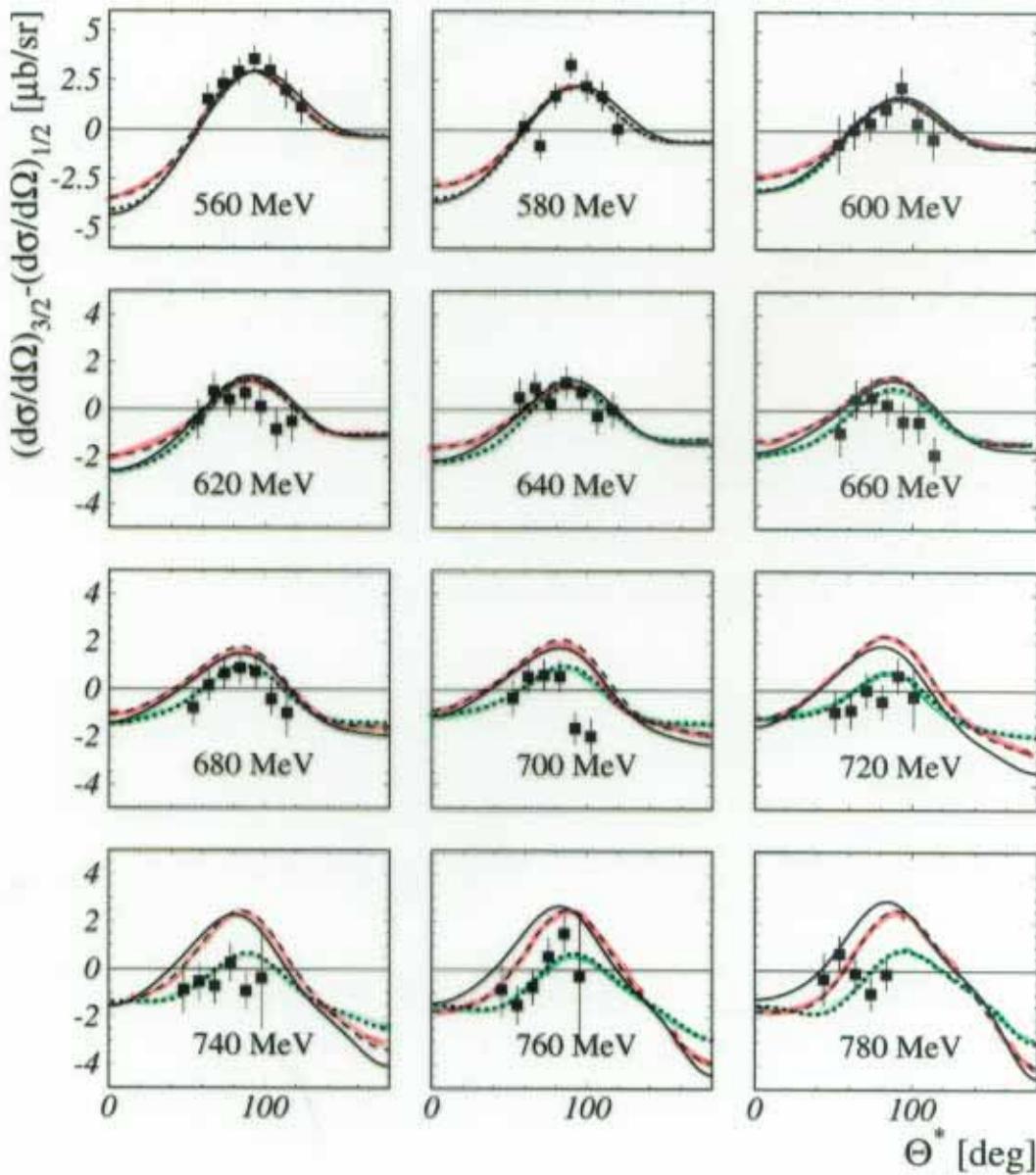


Results



$$\vec{\gamma} \vec{p} \rightarrow p \pi^0$$





■ GDH data: J. Ahrens et al., PRL 88, 232002 (2002)

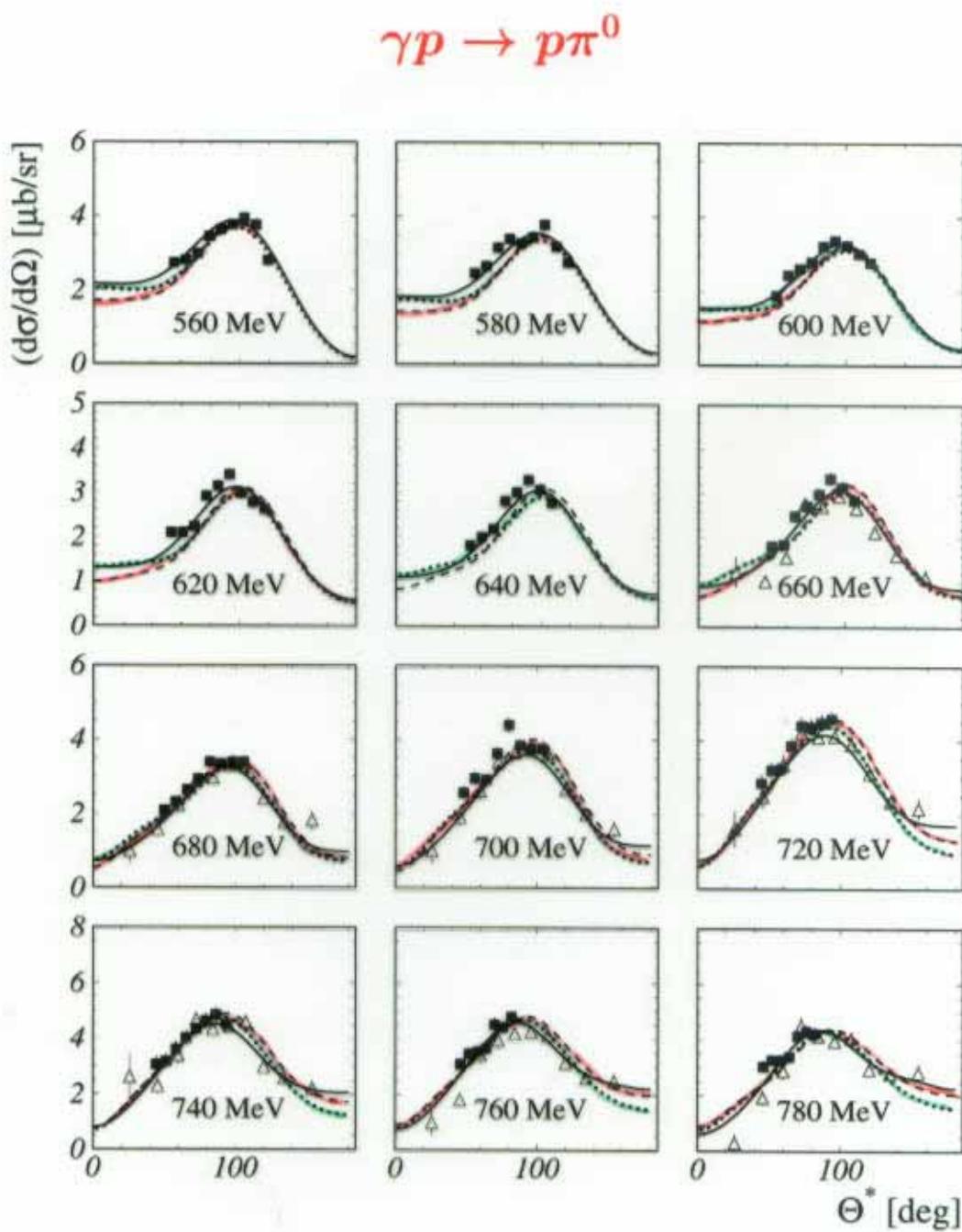
— SAID (σ_{p01})

— UIM

— Modified UIM

$M_{2-}^{1/2} : +11\%$

$E_{2-}^{1/2} : -20\%$



■ GDH data: J. Ahrens et al., PRL 88, 232002 (2002)

△ TAPS data: B. Krusche et al., EPJ A6, 309 (1999)

— SAID (*sρo1*)

— UIM

Modified UIM

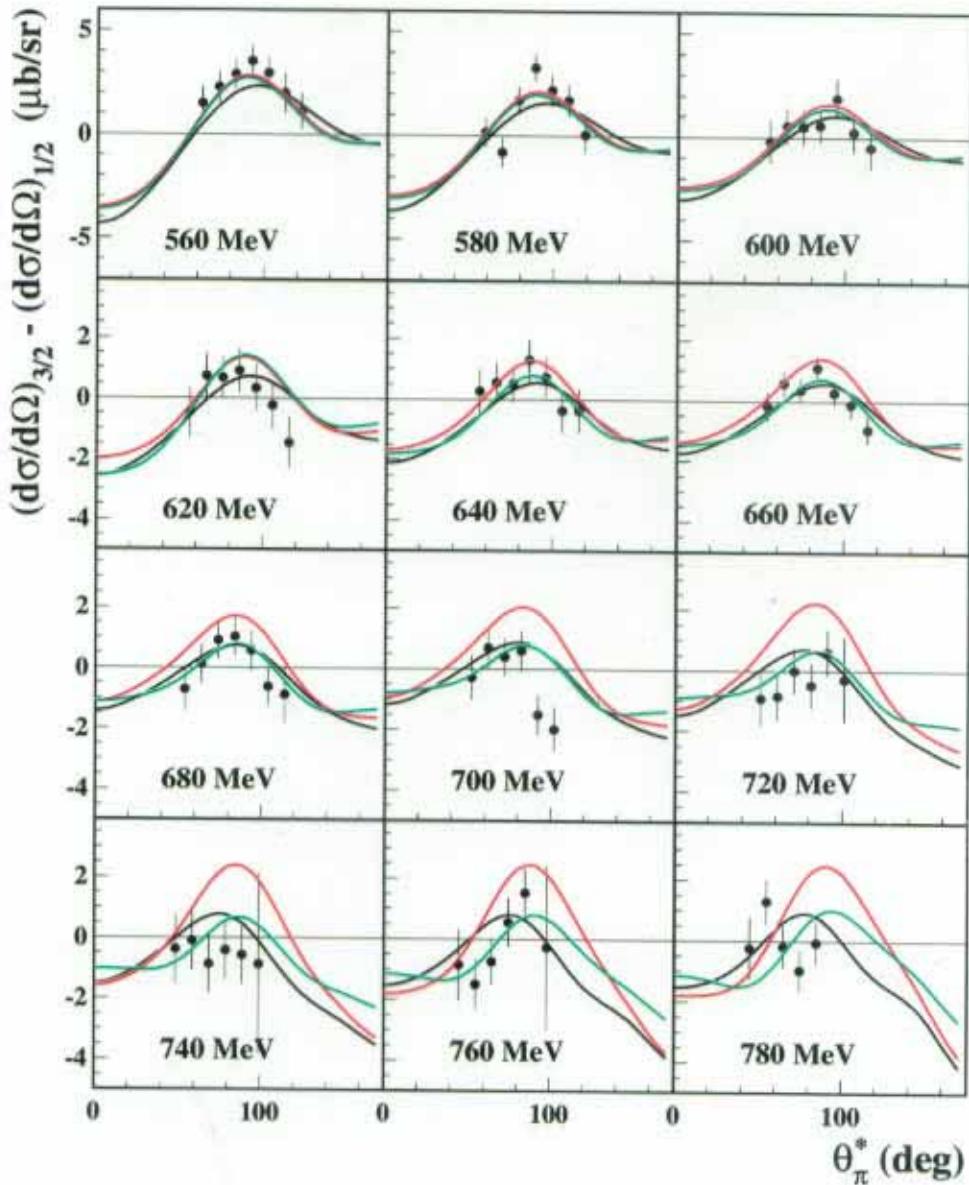
*D*₁₃ Helicity amplitudes

Solution	<i>A</i> _{1/2}	<i>A</i> _{3/2}
Standard UIM	-17	164
Modified UIM	-38 ± 3	147 ± 10
PDG “best estimate”	-24 ± 9	166 ± 5

Units: 10⁻³ GeV^{-1/2} $W_R = 1520$ MeV; $\Gamma_\pi/\Gamma = 0.55$; $\Gamma = 120$ MeV

Solution	<i>A</i> _{1/2}	<i>A</i> _{3/2}
SAID (SM02K)	-24 ± 2	135 ± 2

Units: 10⁻³ GeV^{-1/2} $W_R = 1517$ MeV; $\Gamma_\pi/\Gamma = 0.63$; $\Gamma = 109$ MeV



- GDH data: J. Ahrens et al., PRL 88, 232002 (2002)

— SAID (solution SM02K)

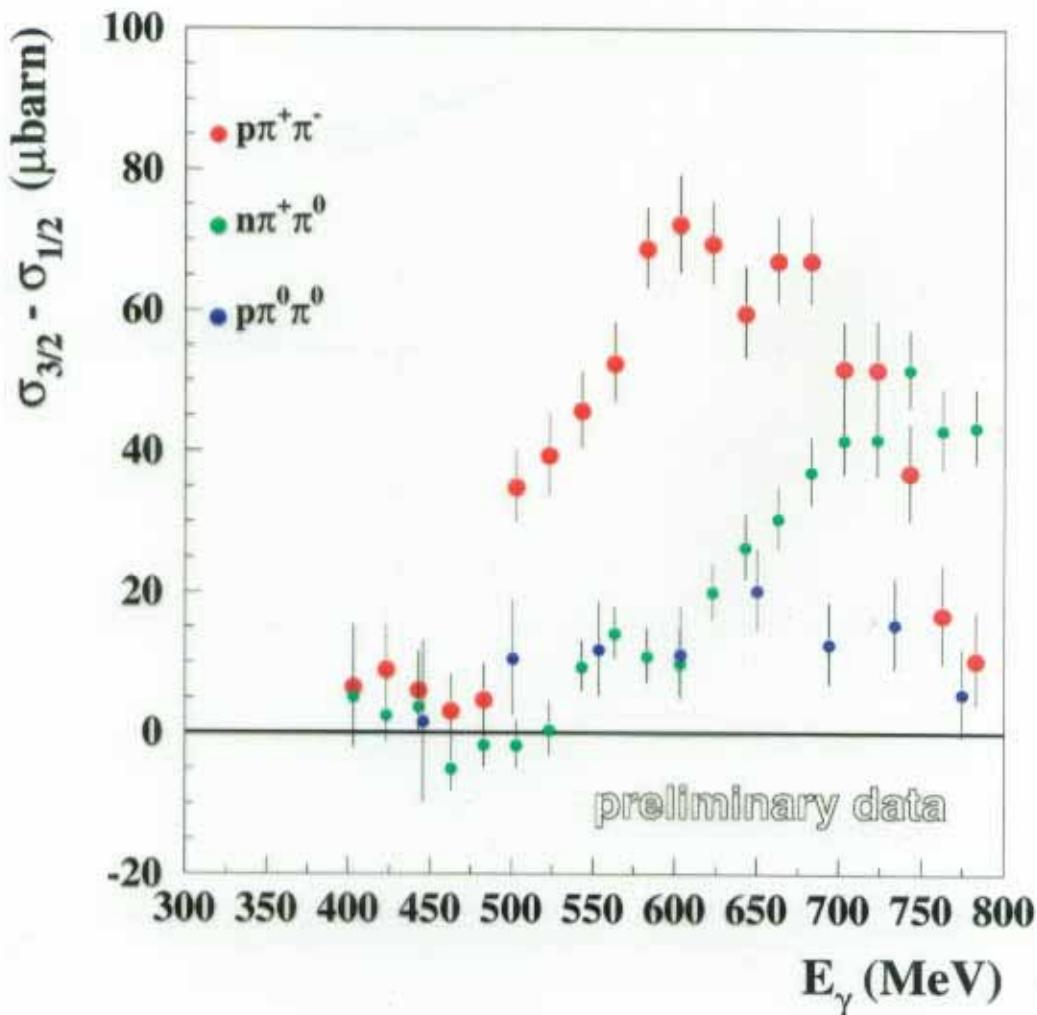
— UIM

— Modified UIM

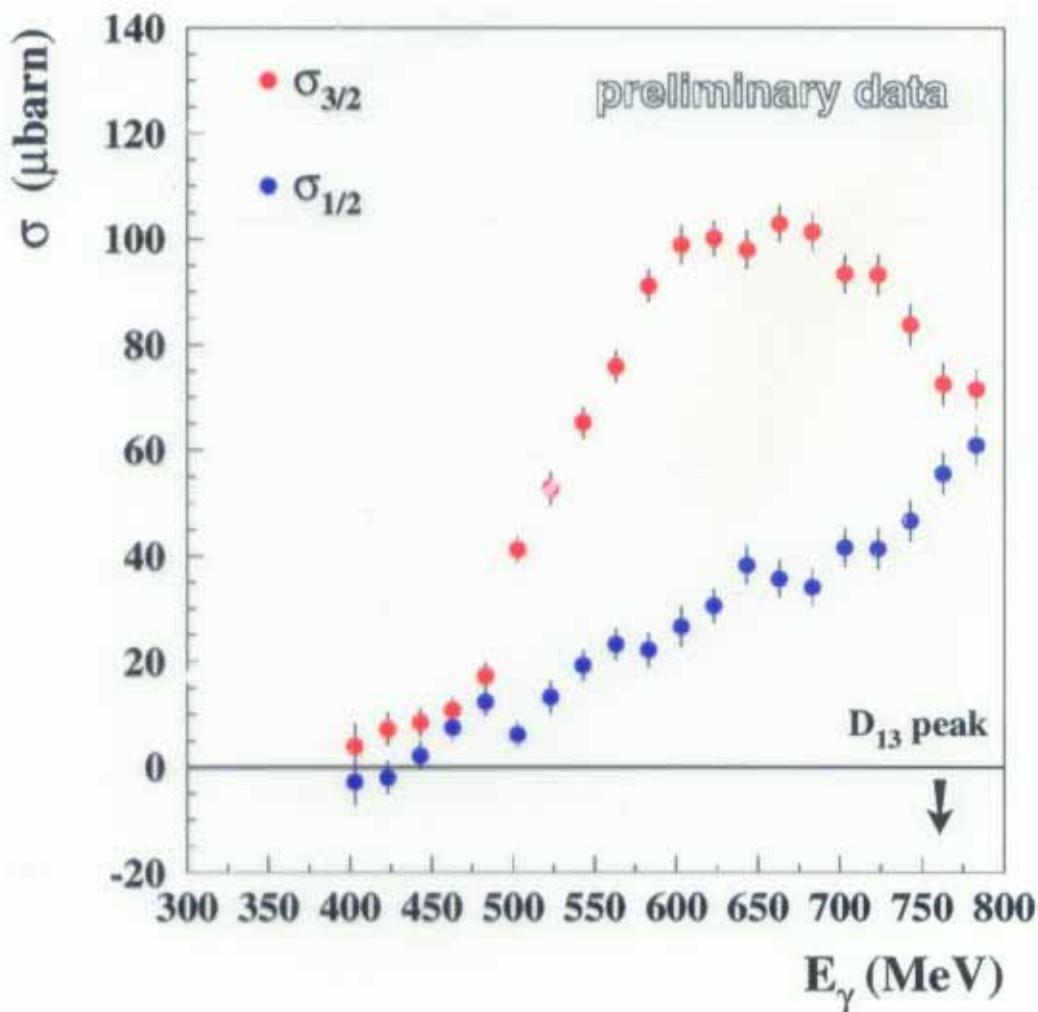
$M_{2-}^{1/2} : +11\%$

$E_{2-}^{1/2} : -20\%$

$$\vec{\gamma} \vec{p} \rightarrow N\pi\pi$$

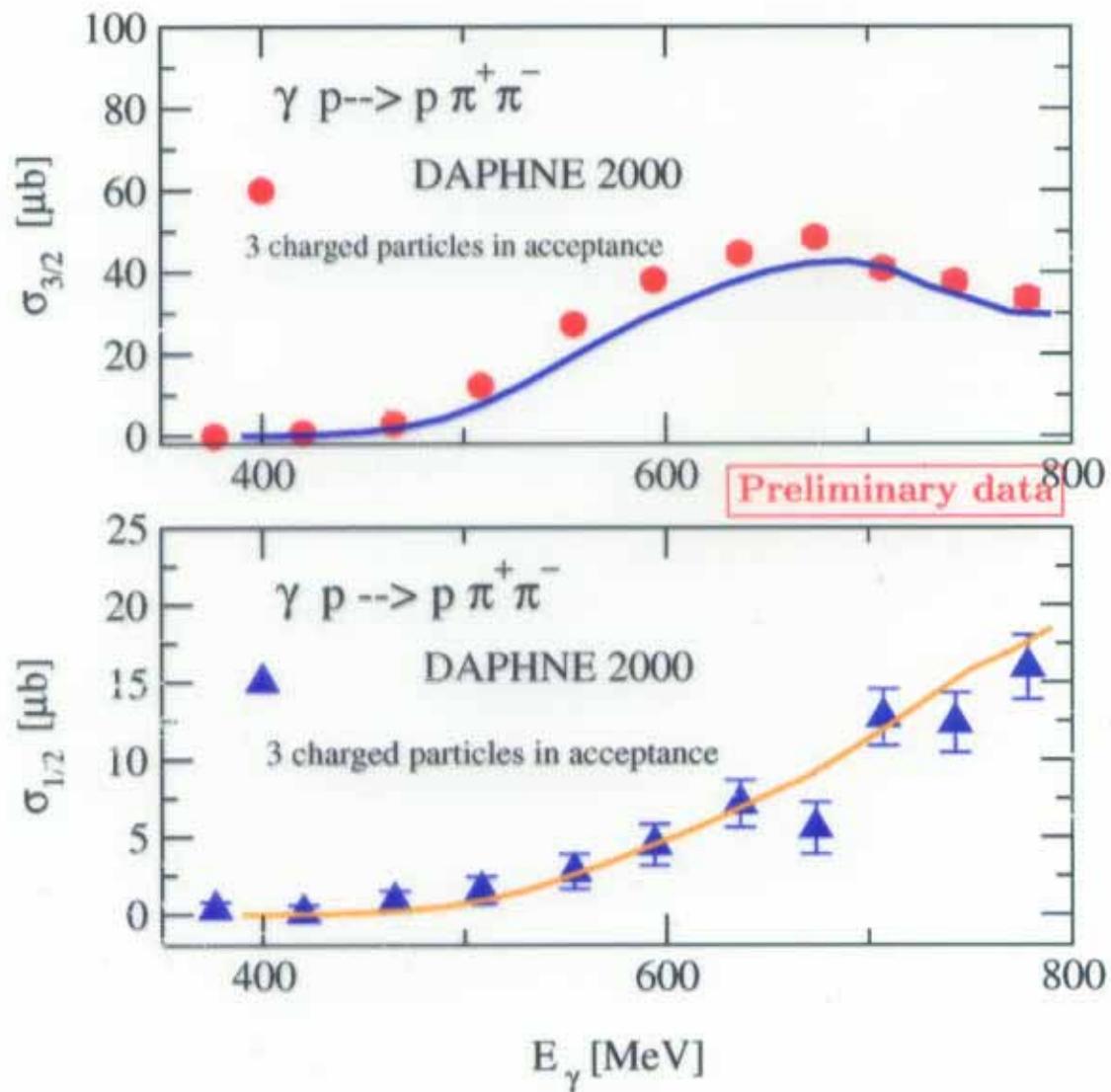


- extrapolation into full acceptance performed using a pure phase space model ($\pi^+\pi^-$; $\pi^0\pi^0$) or Nacher-Oset model ($\pi^+\pi^0$)
- correction is $\simeq 20\%$ of the measured cross section for $n\pi^+\pi^0$, $\simeq 30\%$ for $p\pi^+\pi^-$, and $\simeq 50\%$ for $p\pi^0\pi^0$
- using unpolarized cross sections → separation between $\sigma_{3/2}$ and $\sigma_{1/2}$

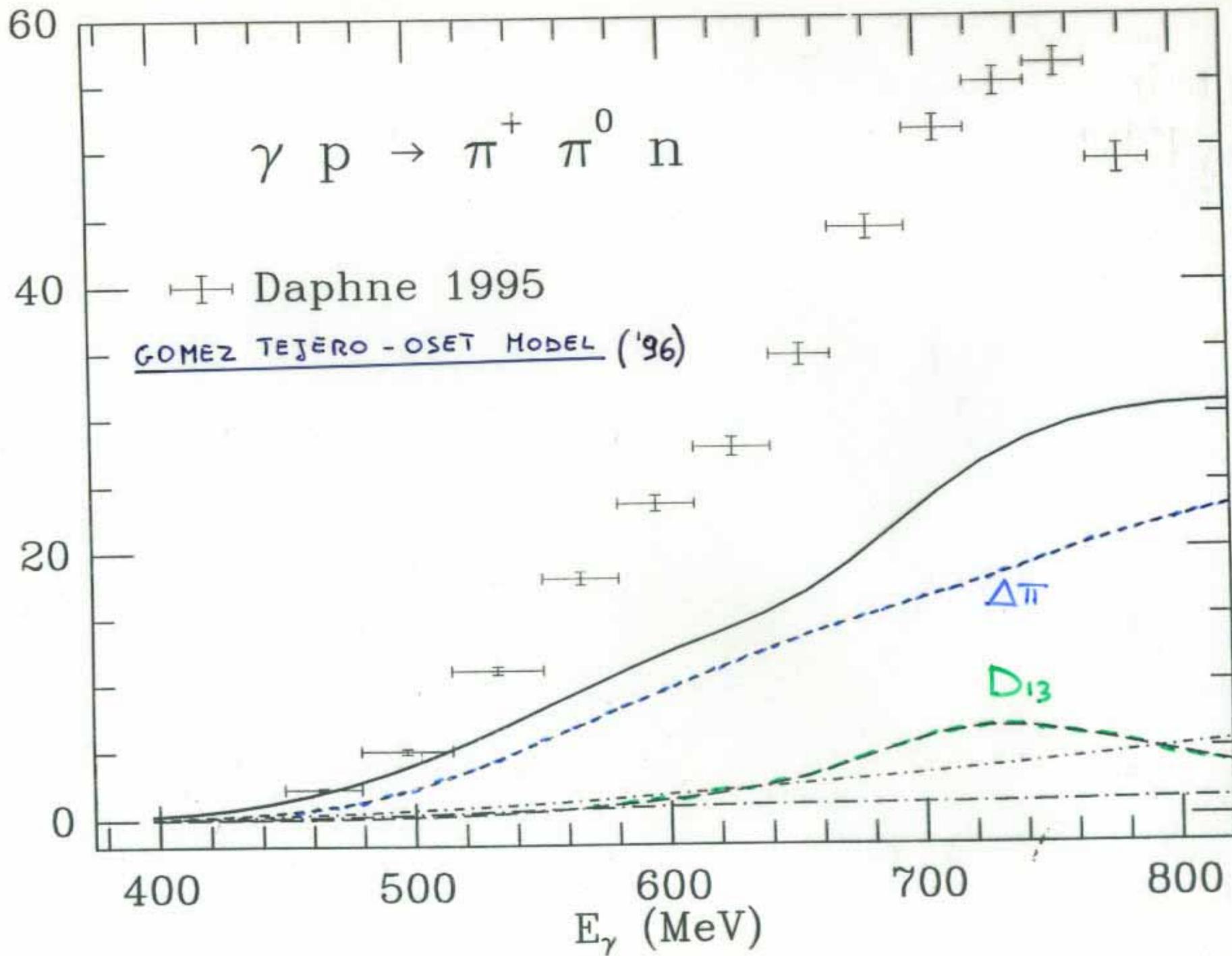


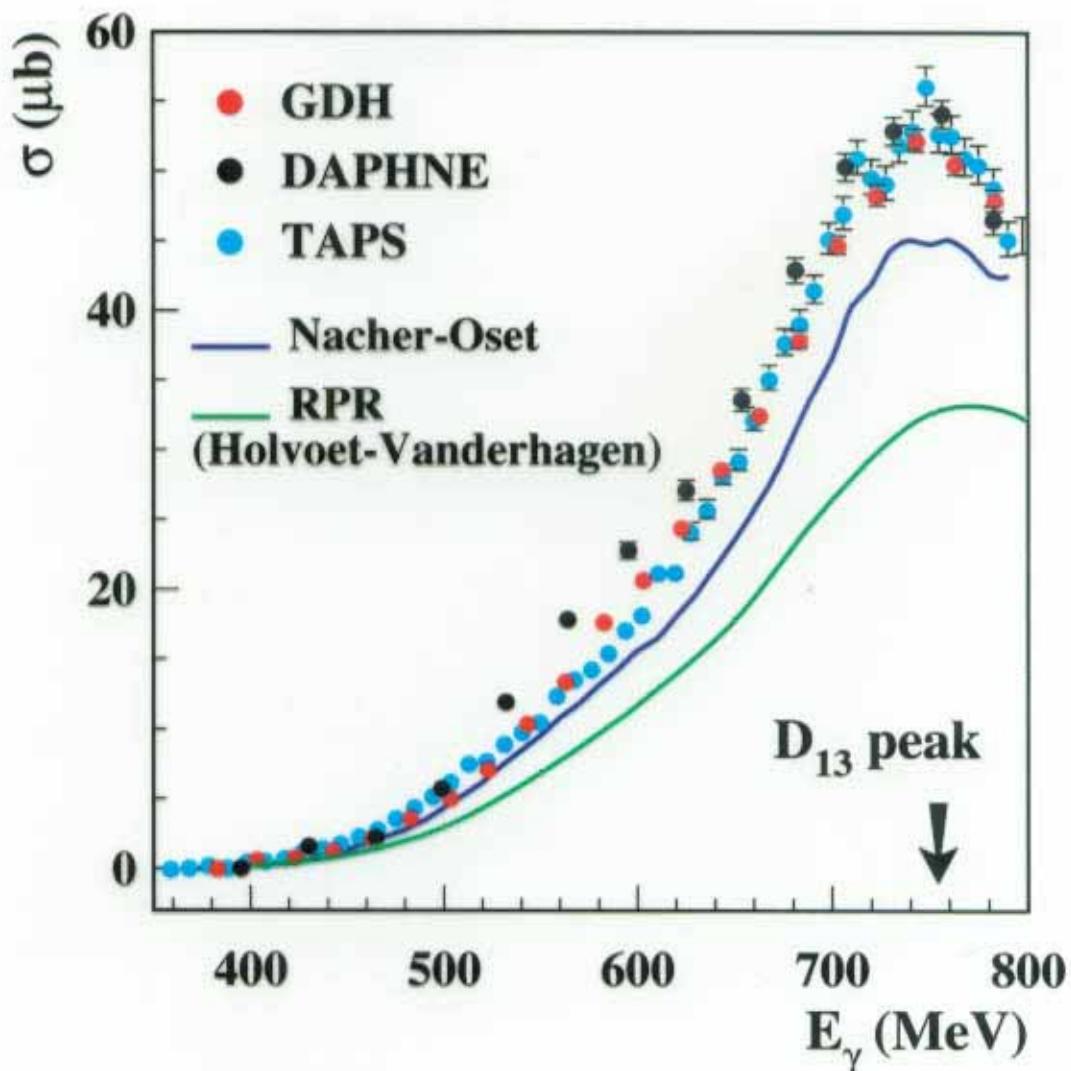
- D_{13} plays a (very) small role
- Most important mechanism: $\gamma p \rightarrow \Delta \pi$
- S -wave $\pi \Rightarrow \sigma_{3/2}$ is dominant

Comparison with Models



- Nacher-Oset model (nucl-th/0106005)
- Data inside DAPHNE acceptance and only for $\simeq 50\%$ of the collected events (\Rightarrow 3 charged particles in DAPHNE)





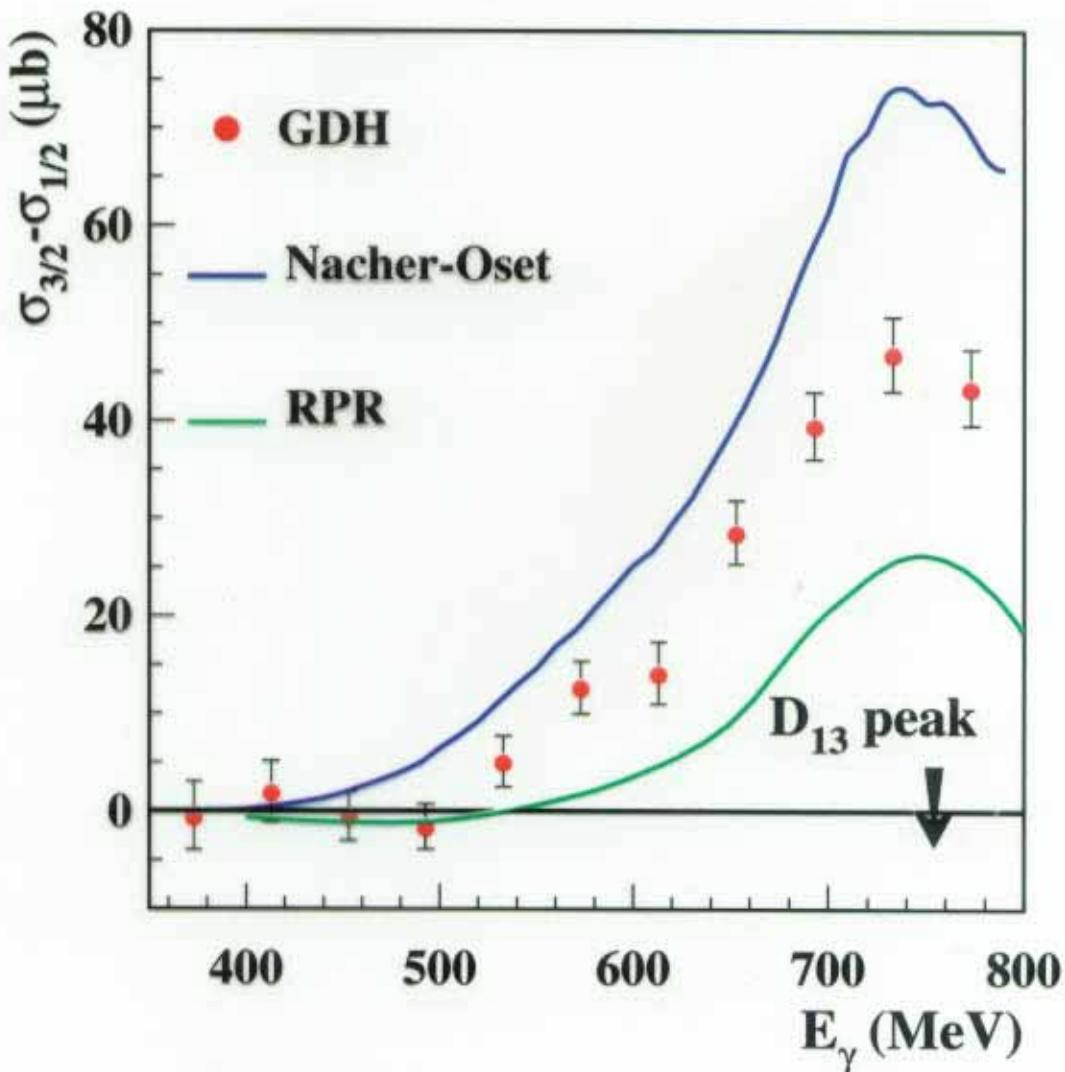
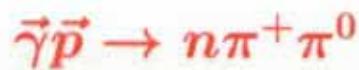
GDH data: J. Ahrens et al., to be submitted

DAPHNE data: A. Braghieri et al., PLB 363, 46 (1995)

TAPS Data: W. Langgartner et al., PRL 87, 052001 (2001)

— J.Nacher, E.Oset, NPA 697, 372 (2002) ($D_{13} \rightarrow n\rho^+$)

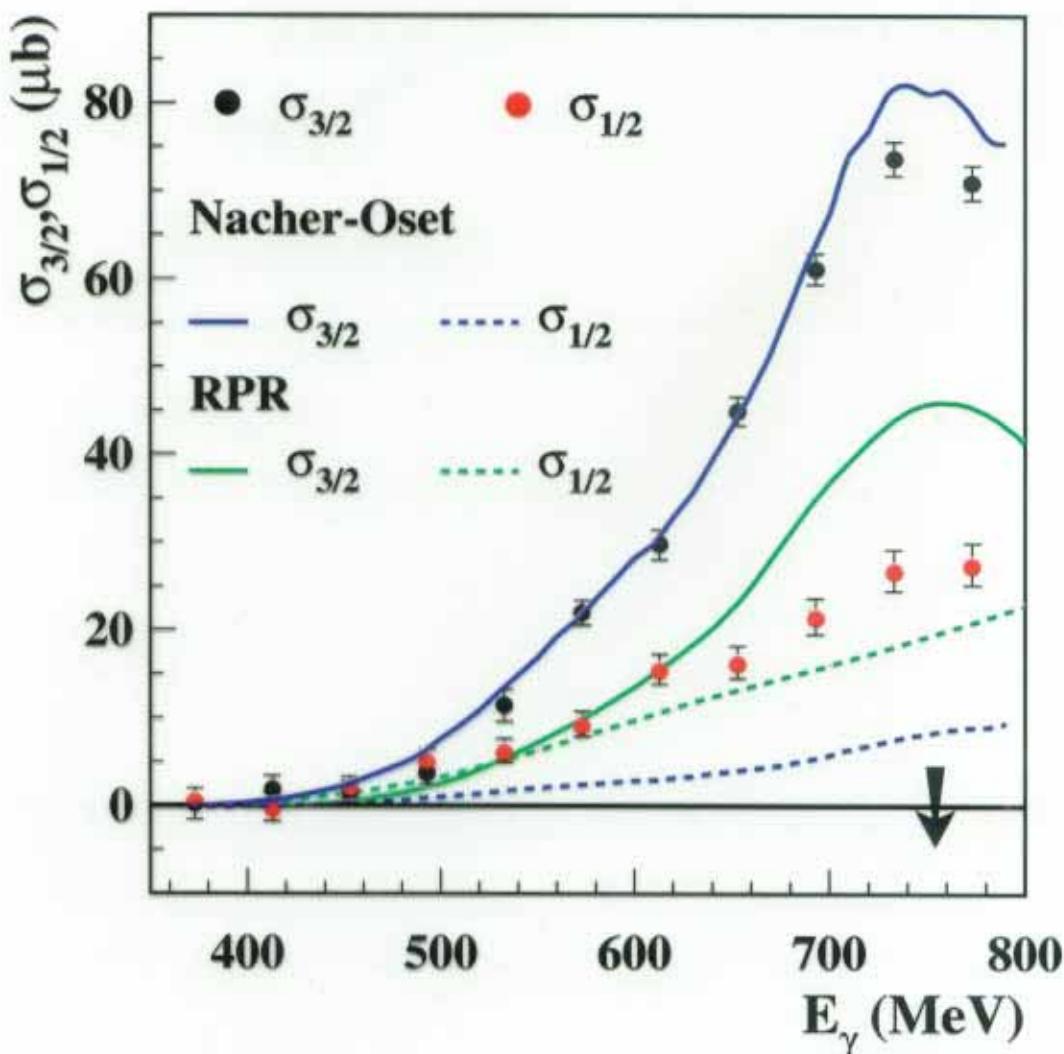
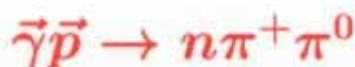
— H.Holvoet, PhD thesis, University of Gent (2001)



GDH data: J. Ahrens et al., to be submitted.

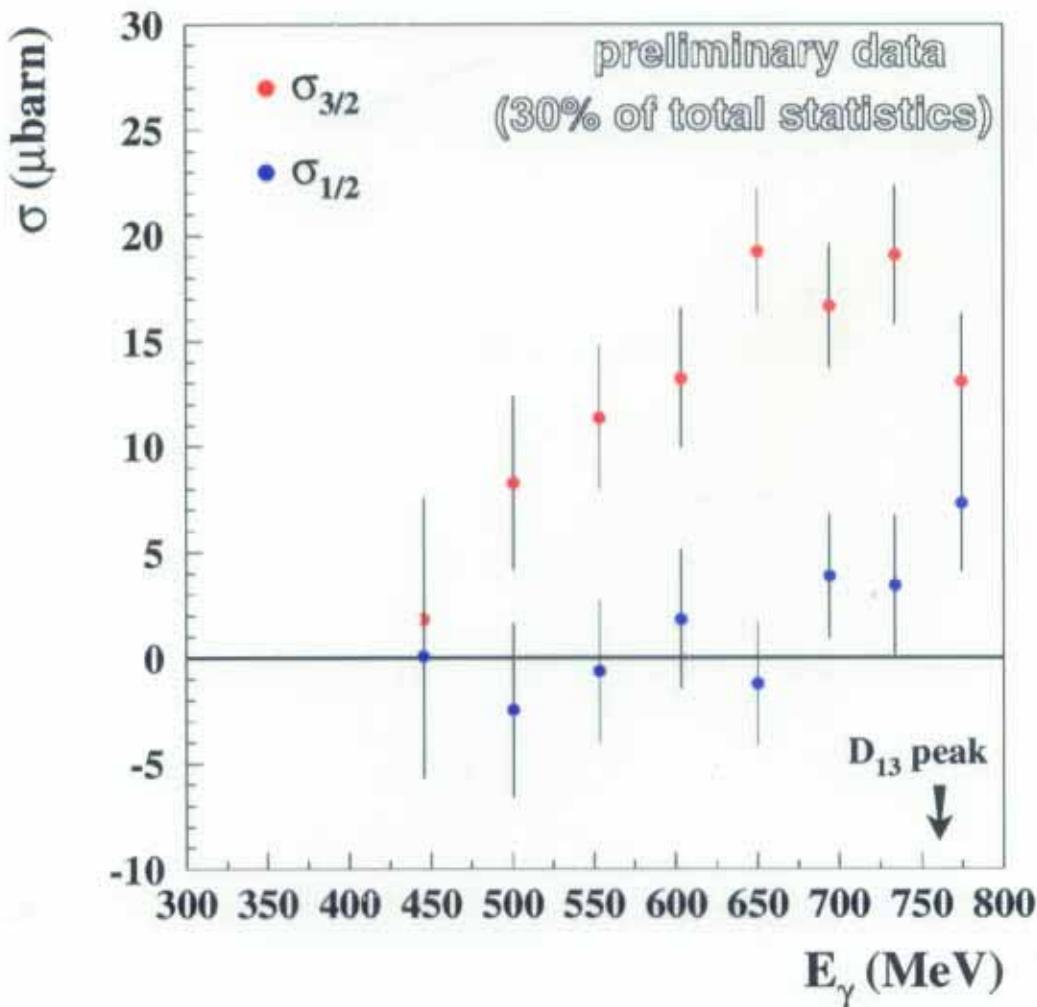
— J.Nacher, E.Oset, NPA 697, 372 (2002) ($D_{13} \rightarrow n\rho^+$)

— H.Holvoet, PhD thesis, University of Gent (2001)



- From previous DAPHNE (PRC 60, 055201 -1999-) and TAPS (PRL 87, 052011 -2001-) data (invariant mass distributions): $\pi\pi$ correlation in the final state
- The most important contribution is given by $\gamma p \rightarrow D_{13} \rightarrow n \rho^+$
- Non-negligible background mechanisms: $\gamma p \rightarrow N \rho (\overline{D} \Delta)$

$$\vec{\gamma} \vec{p} \rightarrow p \pi^0 \pi^0$$



- Statistics too low (and extrapolation too large) to draw “sound” conclusions
- D_{13} plays an important role, as expected from $\pi\pi, p\pi$ invariant mass distributions (TAPS, EPJ A9, 5 -2000-)
- No (relevant) contribution is given by P_{11}

GDH sum rule on the neutron

- Total cross section measurement \Rightarrow “inclusive” method (no partial channel separation)

- ^2H : $\mu \sim \mu_p + \mu_n \Rightarrow \begin{array}{c} \uparrow \uparrow \\ p \ n \end{array}$

$$I_{exp}^d \sim I_{GDH}^n + I_{GDH}^p + I_{nucl.eff.}$$

- ^3He : $\mu \sim \mu_n \Rightarrow \begin{array}{c} \uparrow \downarrow \uparrow \\ p \ p \ n \end{array}$ (S state with $\sim 90\%$ prob.)

$$I_{exp}^{he} \sim I_{GDH}^n + \alpha I_{GDH}^p + I_{nucl.eff.}$$

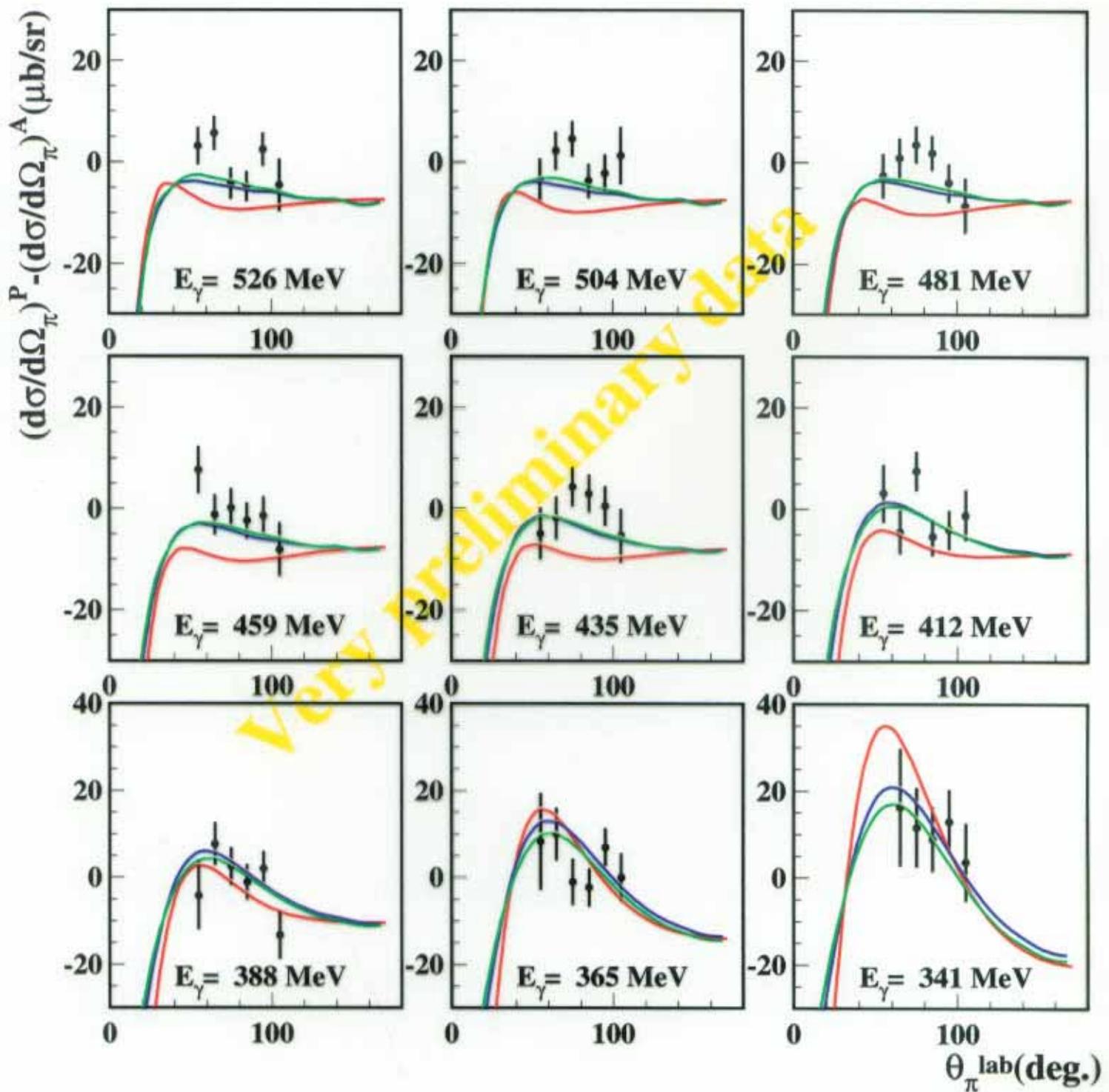
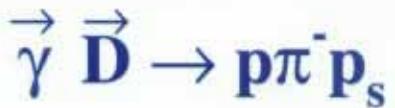
$$\alpha \simeq 0.1$$



smaller corrections

^3He better suited to get an estimate of
 I_{GDH}^n

^2H better suited to measure $\vec{\gamma}\vec{n} \rightarrow N\pi(\pi)$



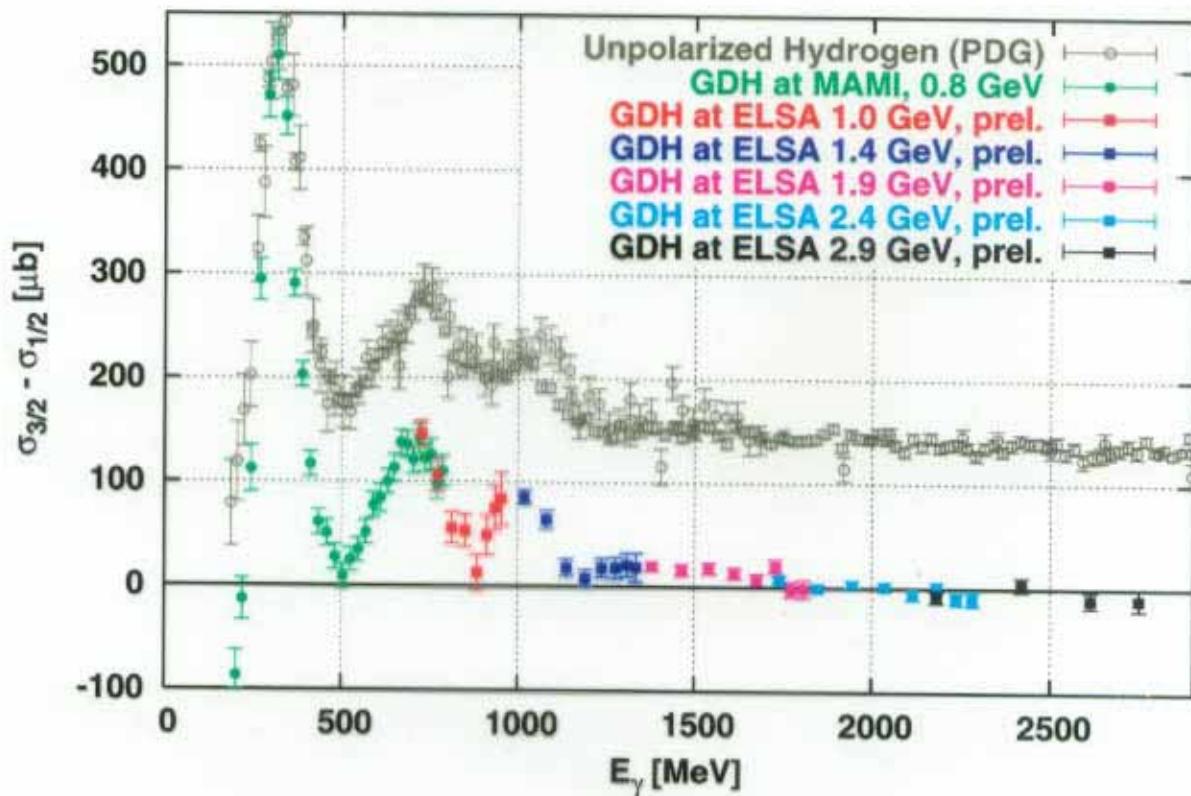
Arenhövel $\gamma n \rightarrow p\pi^-$ free

Arenhövel $\gamma D \rightarrow p_s p\pi^-$ IA

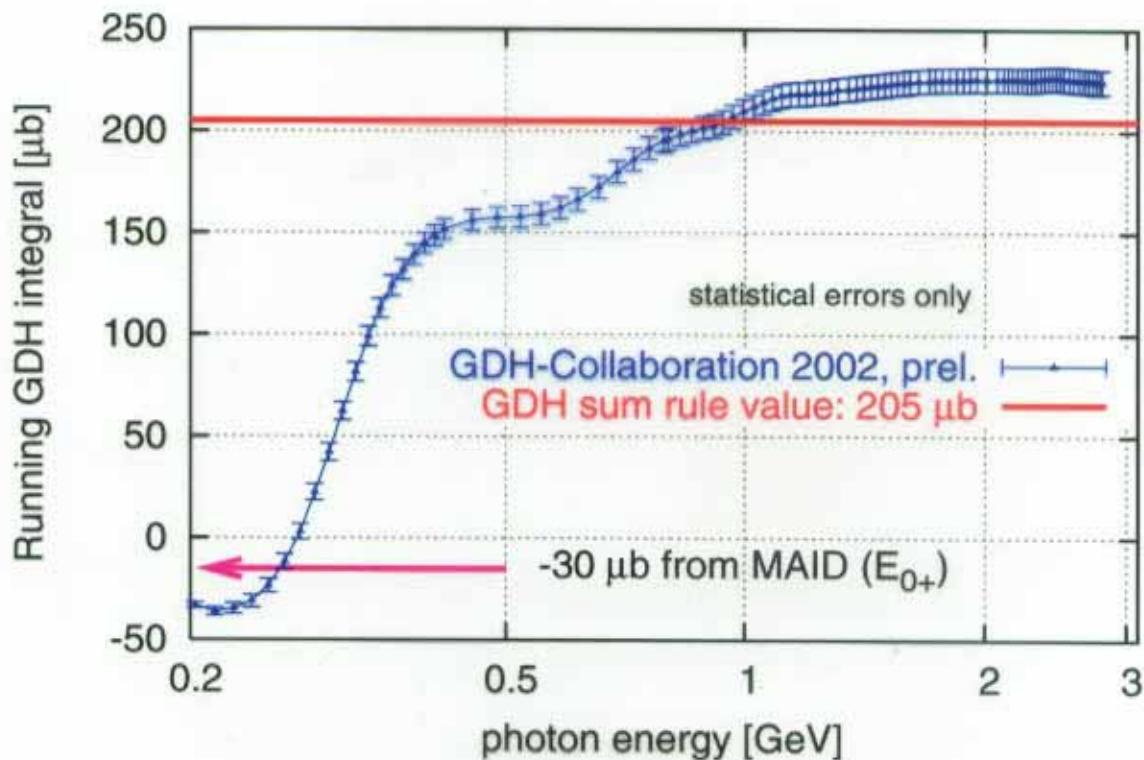
Arenhövel $\gamma D \rightarrow p_s p\pi^-$ IA+FSI

Model: E.Darwish, Ph.D.Thesis, Mainz(02)

Data: C.A.Rovelli, Dipl.Thesis, Pavia(02)

Inclusive $\vec{\gamma}\vec{p} \rightarrow \text{hadrons}$ 

- Mainz (MAMI) data: J. Ahrens et al. PRL 87, 02203 (2001)
- Bonn (ELSA) data: Preliminary

Inclusive $\vec{\gamma}\vec{p} \rightarrow \text{hadrons}$ 

- Running GDH integral = $\int_{200 \text{ MeV}}^{E_\gamma^{\max}} \frac{(\sigma_{3/2} - \sigma_{1/2})}{E_\gamma} dE_\gamma$
- $I_{GDH}[140 - 200 \text{ MeV}] \simeq -30 \mu\text{b}$ (multipole analyses)
- $I_{GDH}[> 3 \text{ GeV}] \approx [-10, -20] \mu\text{b}$ (Regge-like models)

Final remarks

- After a long “hunt” for the GDH sum rule, we are almost there.
The (generalized) GDH sum rule is on the schedule of all existing $e/\mu/\gamma$ facilities
- Results from Mainz/Bonn:
 - the GDH sum rule value for the proton is (slightly) over-shot at $E_\gamma = 3 \text{ GeV}$
If I would gamble, I would bet that the GDH sum rule (on the proton) is verified.
 - Spin polarizability $\gamma_0 \sim -80/-90 \cdot 10^{-6} \text{ fm}^4$
- Helicity dependence of the $N\pi(\pi)$ channels is a powerful tool to better understand the baryon resonance properties
 $\Rightarrow D_{13}$ properties are different from the “expected” ones
- Additional measurements both on proton and on the neutron (higher photon energies; all relevant partial channels) are needed to get a more complete picture both on the resonance properties and on the GDH sum rule.