

## Survey of Lattice Results for Baryon States.

David Richards (LHPC and Jefferson Laboratory)

- Lattice QCD spectrum recipe
- Quenched approximation
- The quenched baryon spectrum - so far
- Full QCD, light quarks. . .
- Decays and transitions
- Further and outlook. . .

In principle, the recipe is to determine the mass of particle  $P$  is straightforward:

- Choose an *interpolating operator*  $\mathcal{O}_P$  that has large overlap:

$$\langle 0 | \mathcal{O}_P | P \rangle \neq 0$$

- Construct the time-sliced correlator

$$C(t) = \sum_{\vec{x}} \langle \mathcal{O}(\vec{x}, t) \mathcal{O}^\dagger(\vec{0}, 0) \rangle$$

- Insert a complete set of states

$$\begin{aligned}
 C(t) &= \sum_{\vec{x}} \sum_P \int \frac{d^3k}{(2\pi)^3 2E(\vec{k})} \langle 0 | \mathcal{O}(\vec{x}, t) | P(\vec{k}) \rangle \langle P(\vec{k}) | \mathcal{O}^\dagger(\vec{0}, 0) | 0 \rangle \\
 &\longrightarrow \sum_P \frac{|\langle 0 | \mathcal{O} | P \rangle|^2}{2m_P} e^{iM_P t} \longrightarrow \sum_P \frac{|\langle 0 | \mathcal{O} | P \rangle|^2}{2m_P} e^{-M_P t} \text{ Euclidean}
 \end{aligned}$$

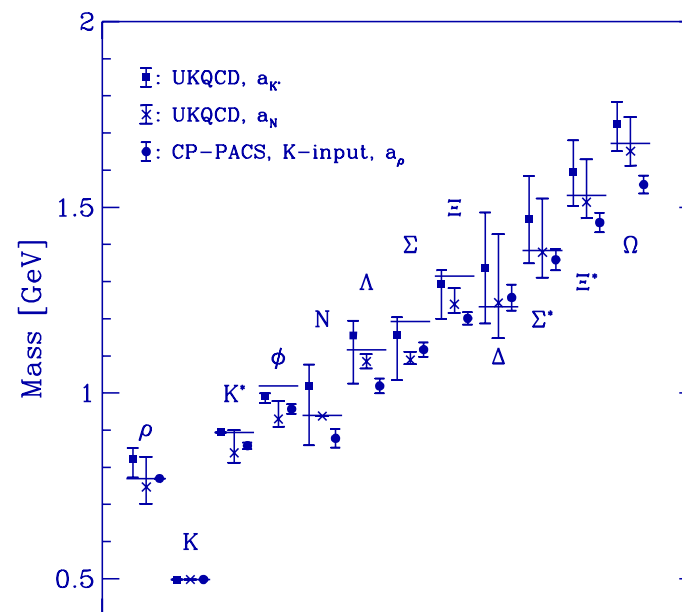
- **time-slicing** puts the intermediate states at rest.
- At large times, correlator dominated by **lightest state**

## Hadron Spectrum - Benchmark of LQCD

It is the most precise and extensively pursued lattice calculation.

### Final Quenched Spectrum

UKQCD, PRD62 (2000), 054506



Quenched Spectrum Agrees with  
Experiment to 10%

Much work behind the scenes...

- Continuum extrapolation  $a \rightarrow 0$
- Extrapolation  $V \rightarrow \infty$
- Chiral extrapolation  $M_{PS} \rightarrow M_{\pi}$

*Inconsistency in meson sector resolved in full QCD*

## Baryon Operators

- Flavour structure
- Parity
- Angular Momentum

Cubic group admits only three irreducible ray representations

### Rep. Spin cpts.

$G_1$	$1/2, 7/2, \dots$
$H$	$3/2, 5/2, 7/2, \dots$
$G_2$	$5/2, 7/2, \dots$

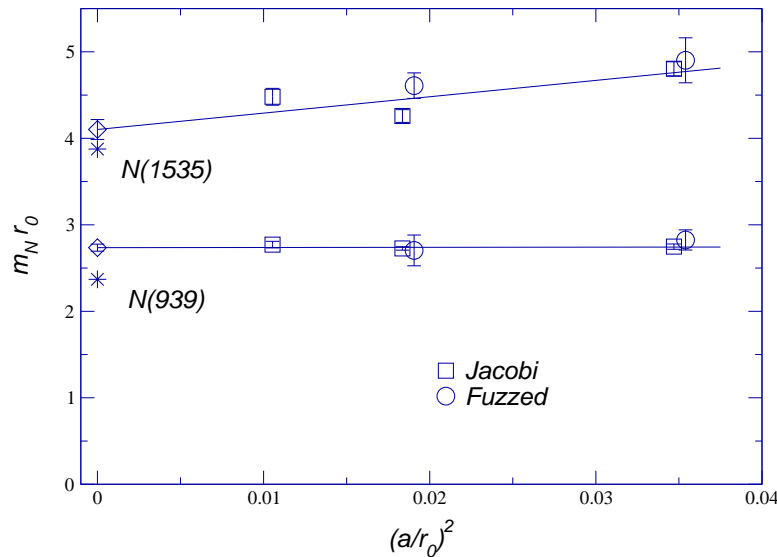
Measure local interpolating operators for the nucleon

$$\begin{cases} N_1^{1/2+} & = \epsilon_{ijk}(u_i^T C \gamma_5 d_j) u_k & \text{Nucleon} \\ N_2^{1/2+} & = \epsilon_{ijk}(u_i^T C d_j) \gamma_5 u_k & \text{Roper} \end{cases}$$

$N_1(N_2)$  connects upper (lower) spinor components in diquark piece -  $N_2$  vanishes in NR limit.

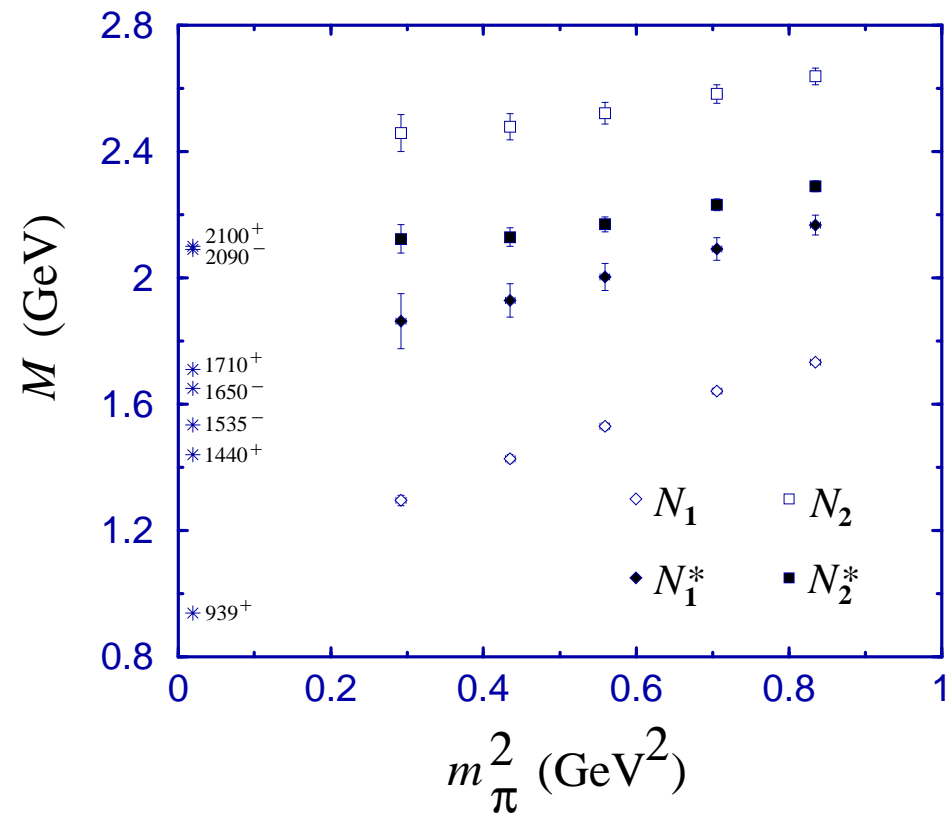
Flurry of lattice activity aimed at extracting lightest states of both parities in each channel:-

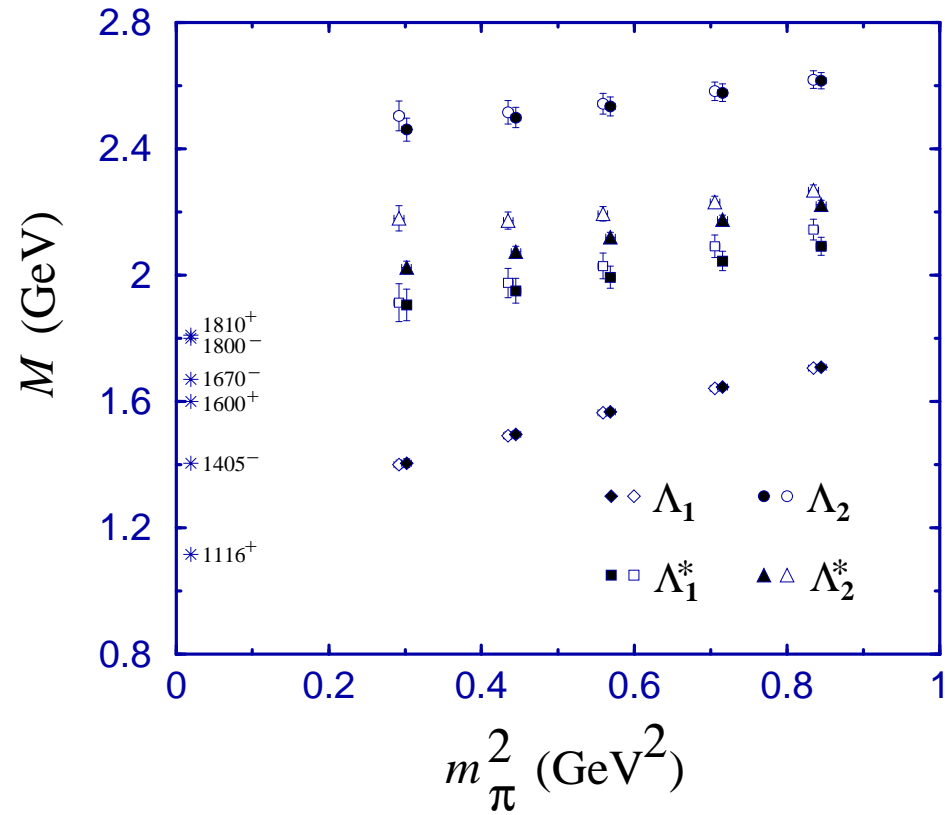
- $D_{234}$  action - Frank Lee and Derek Leinweber
- Domain-wall fermions - BNL/Riken
- Clover - LHPC/QCDSF/UKQCD; anisotropic clover - LHPC
- Improved gauge/FLIC actions - Adelaide Group



LHPC/QCDSF/UKQCD

FLIC action, Melnitchouk *et al*





All calculations quenched, at large pion mass, use simple, three-quark operators.

Ordering of states  $\leftrightarrow$  quark model

$$N^{1/2+} < N^{1/2-} < N^{1/2+}.$$



## Roper resonance

- Is Roper lurking as excited state in  $N_1$  correlator - different interpolating operator? - matrix analysis suggests not (BNL, PRD 65, 074503).
- Full QCD?
- Light pion mass?

## Bayesian analysis of correlators:

- Wilson - Sasaki *et al* - hep-lat/0209059
- Clover/dynamical - C. Marnard, DGR - hep-lat/0209165
- Overlap fermions - - Kentucky, F.X. Lee *et al*, hep-lat/0208070

c. 1900 dentists discover Nitrous Oxide - **painless tooth extraction.**

c. 2000 lattice theorists discover Bayesian Statistics - **painless mass extraction**

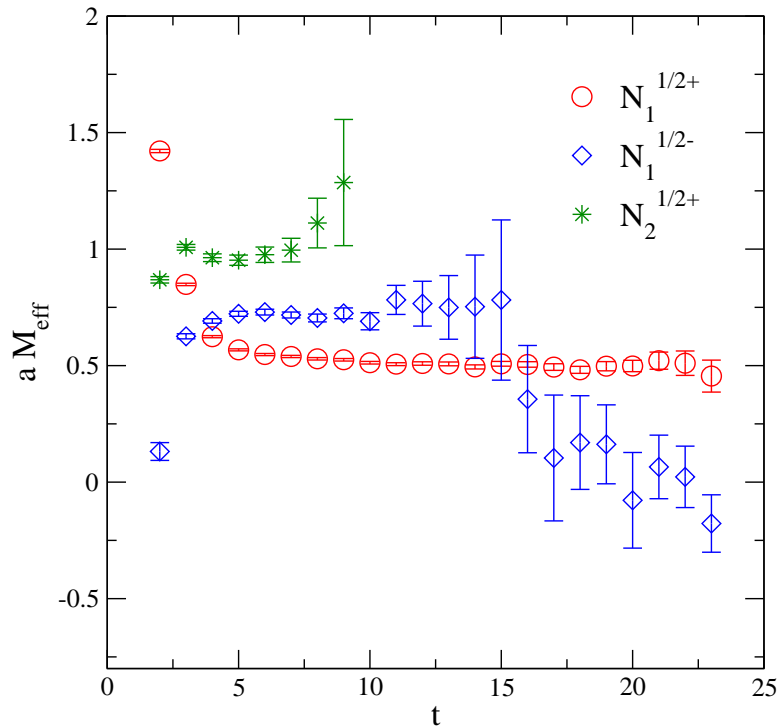
Bayes's Theorem

$$P(M|D \cap I) = \frac{P(D|M \cap I)P(M|I)}{\int dM P(D|M \cap I)P(M|I)}$$

- D is data
- M is model
- I is our prior knowledge

Why?

- Can we obtain more information from the lattice simulations?
- How do we incorporate prior knowledge, e.g.  $m > 0$



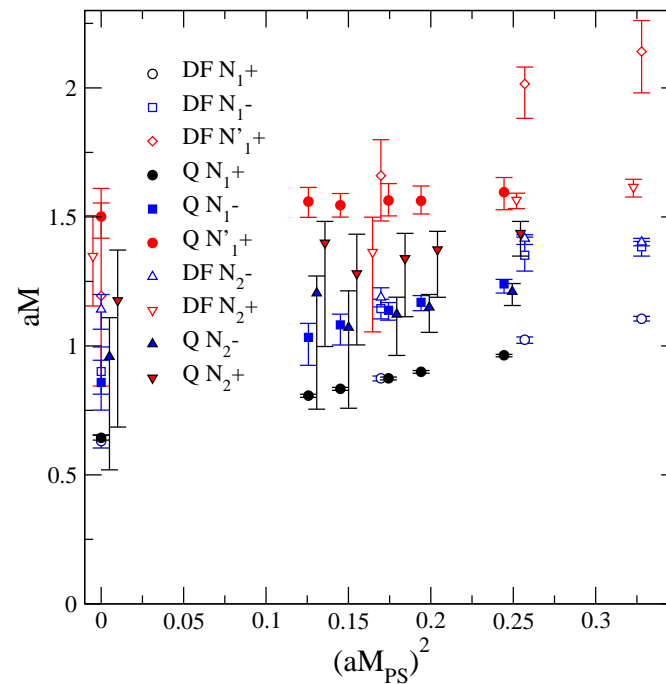
$$m_{\text{eff}} = \ln C(t)/C(t + 1)$$

### Bayesian Priors

What prior knowledge can we incorporate?

- Masses are positive
- Masses are ordered:
- Splitting?

## Comparison of ordering between full and quenched QCD - C. Maynard, DGR (LHPC/UKQCD)



Observed Nucleon Spectrum has ordering  $M^+ < M^- < M'^+$  for pion masses larger than around  $500\text{MeV}$  in full and quenched QCD.

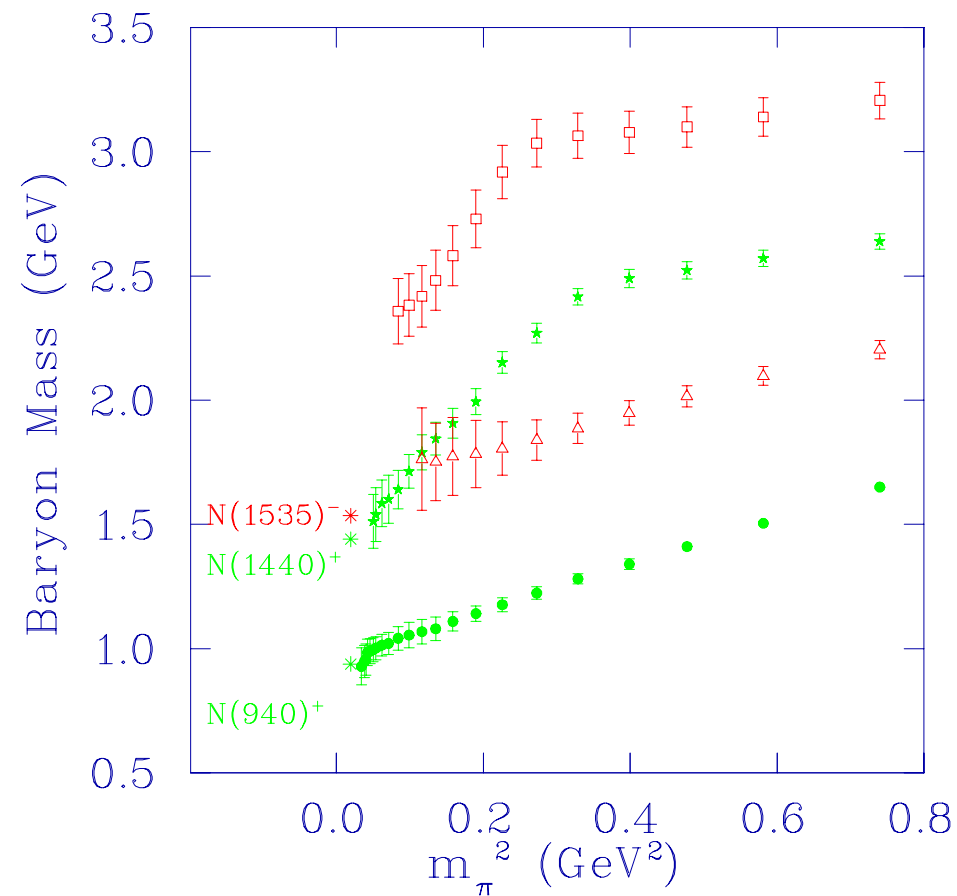
## Extension to light pion masses

- Wilson fermions limited to  $m_\pi \simeq 400\text{MeV}$
- “Smoothed” actions to  $m_\pi \simeq 300\text{MeV}$  - see **FLIC** action.
- Chiral symmetry restored only in **continuum limit**.

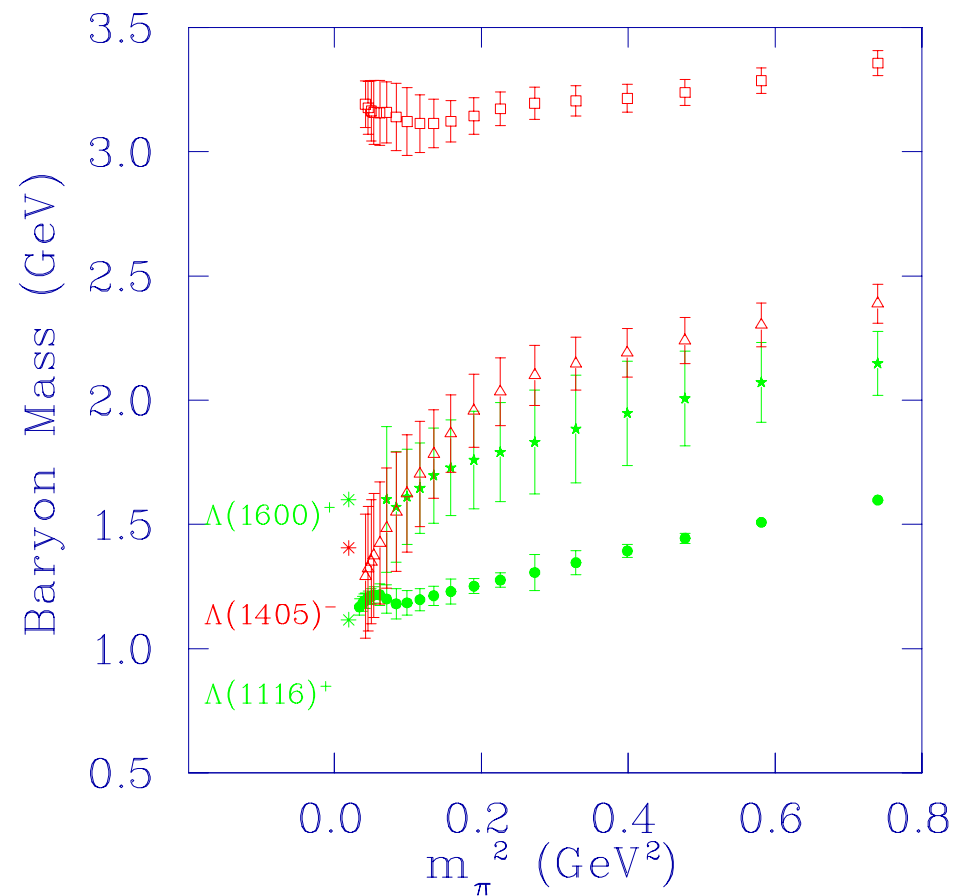
## Overlap or Domain-wall Fermions - Neuberger, Narayanan, Ginsparg/Wilson

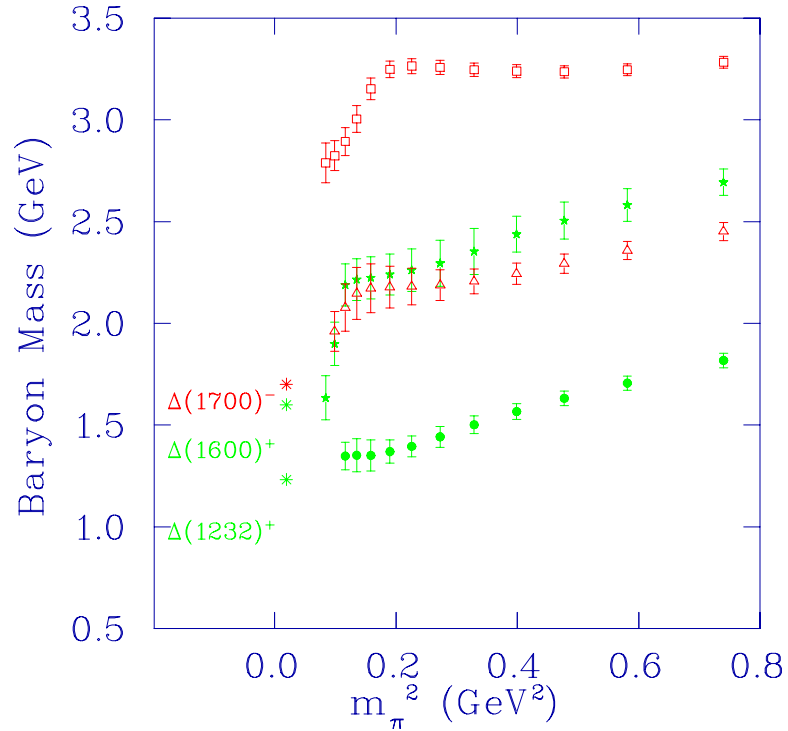
- Possess exact analogue of chiral symmetry at finite lattice spacing.
- No additive mass renormalisation
- Improved operator mixing
- Automatically  $\mathcal{O}(a)$ -improved.

**BUT** ... around 30 times as computationally demanding...

F.X. Lee *et al.*

F.X. Lee *et al.*





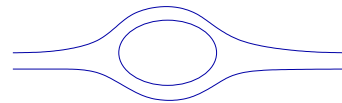
Dramatic reordering of spectrum at light pion mass.

Physics at physical pion mass different to that at heavy quark masses - chiral behaviour/importance of non-analytic terms.



## Decays and Mixing

Great myths of lattice QCD - “Particles don’t decay in the quenched approximation”



(a)



(b)

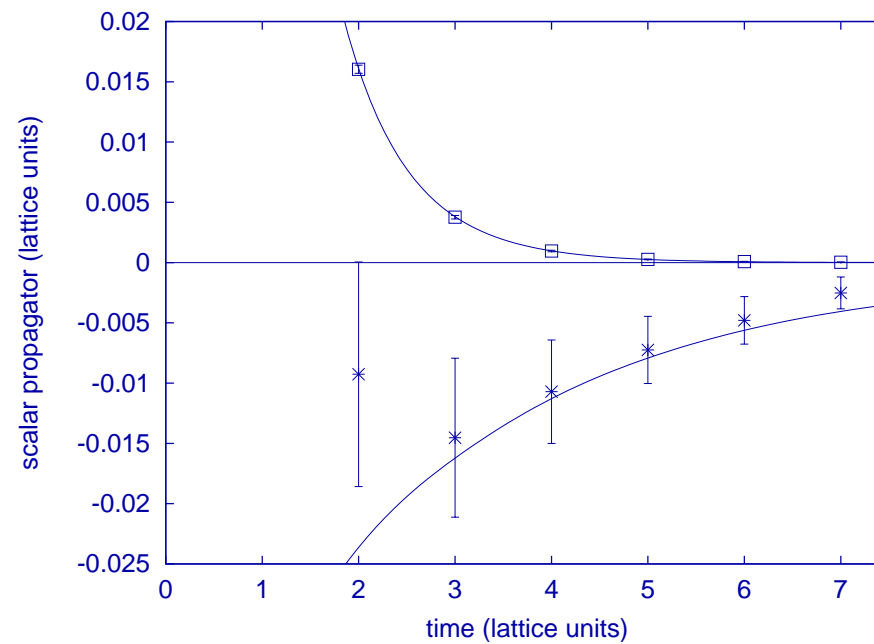


(c)

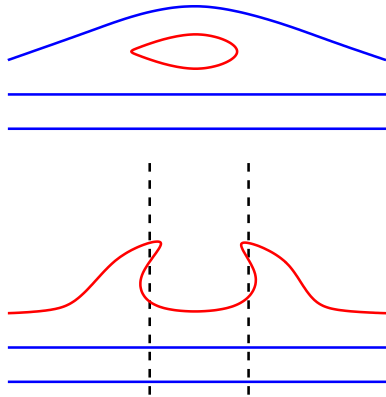
Contributions to the  $\eta'/\pi$  intermediate state in scalar isovector propagator.

Manifest as **non-unitary behaviour** in the hadron correlator.

Bardeen *et al*, PRD 65, 014509 (2001)



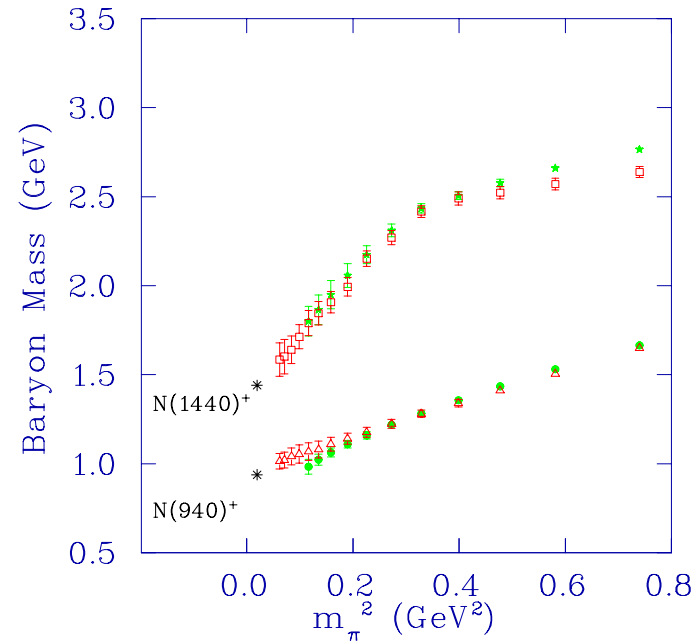
Behaviour well-described by quenched chiral perturbation theory.



Volume-behaviour of two-particle spectrum?

- $m_\pi L \simeq 3$
- Two-particle final states  $\leftrightarrow$  strong  $L$  dependence

F.X. Lee *et al.* (preliminary)



## Beyond the spectrum...

Transition form factors are straightforwardly accessible to lattice calculation  
- Gunnar Bali.

Decays  $A \rightarrow B + C$  complicated because phase information is obscured in Euclidean space - large time correlators dominated by lightest two-body state with minimum momentum - Maiani-Testa Theorem.

Lüscher relates shift in energies of two-particle system in finite box to extract phase-shifts.

Momenta on torus quantised:  $\vec{q} = 2\pi n/L$

Thus, for two pions, say, total  $\pi$ - $\pi$  energy

$$W = 2\sqrt{m_\pi^2 + p_i^2}$$

has discrete values.

Discrete energy eigenvalues of two-pion system shifted by finite size effect.

Lüscher relates the shifts of discrete two-particle energies to the infinite-volume phase shifts.

For zero-momentum state, we have

$$\Delta E = -\frac{2\pi(m_1 + m_2)a_0}{m_1 m_2 L^3} \left[ 1 - c_1 \frac{a_0}{L} + c_2 \left( \frac{a_0}{L} \right)^2 \right]$$

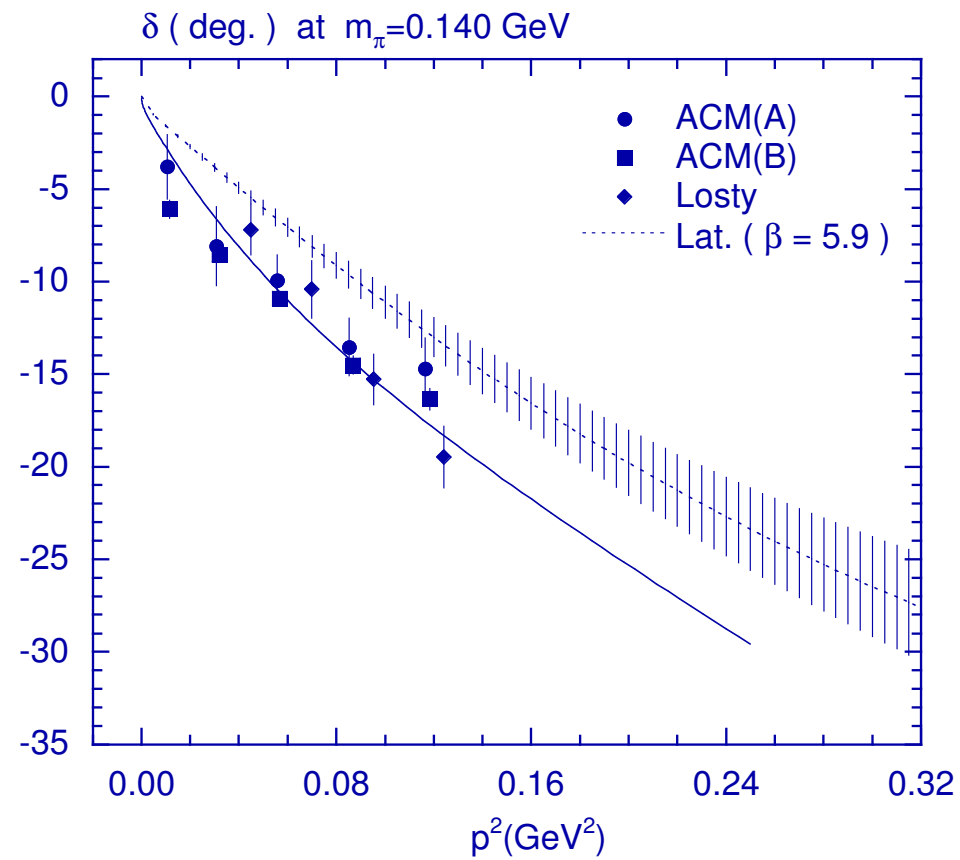
where  $a_0$  is  $S$ -wave scattering length.

Need  $a_0/L$  small for this expansion to be useful.

The method require the extraction of energies from four-point function:

$$C_{h_1 h_2}(t_1, t_2, t_3, t_4) = \sum_{\vec{x}_4} \mathcal{O}_{h_2}(\vec{x}_4, t_4) \sum_{\vec{x}_3} \mathcal{O}_{h_1}(\vec{x}_3, t_3) \sum_{\vec{x}_2} \mathcal{O}_{h_2}^\dagger(\vec{x}_2, t_2) \sum_{\vec{x}_4} \mathcal{O}_{h_2}^\dagger(\vec{x}_4, t_4)$$

$I = 2$  Pion scattering phase shift - Aoki *et al*, hep-lat/0209124



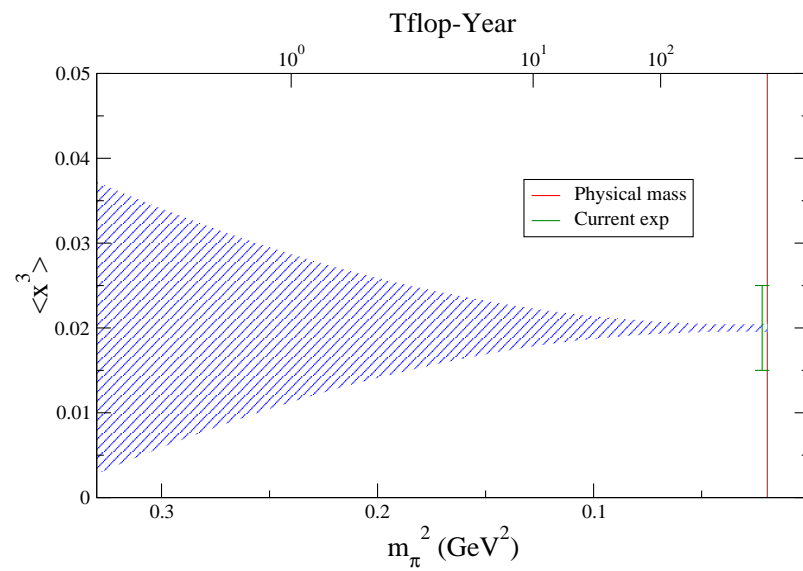
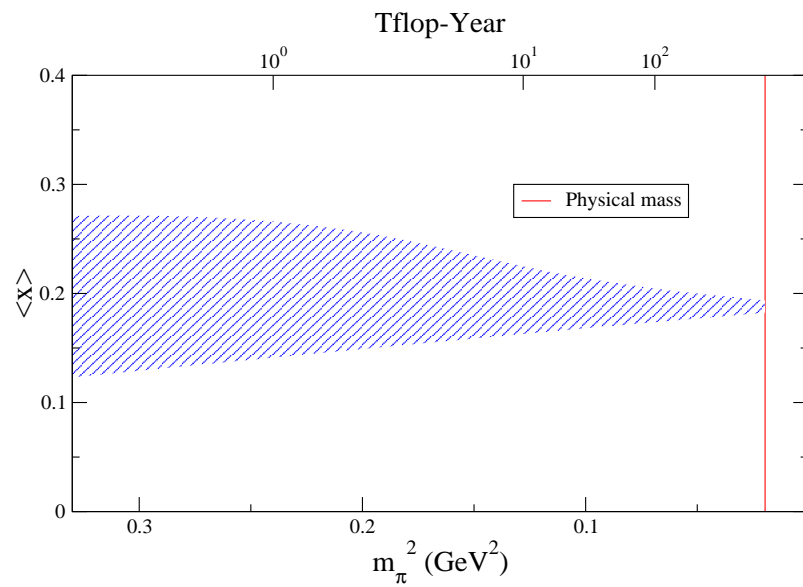
## Summary

- Lattice baryon resonances exciting new application of lattice QCD
- At large masses, resonance spectrum described by quark model.
- Physics is different at light pion mass - domain-wall/chiral fermions.
- Lattice calculations must be matched by careful (quenched) chiral perturbation theory calculations.
- Need to understand finite-volume effects.
- Operators sensitive to “molecules” and excited glue - Feed back with models.
- Many operators  $\leftrightarrow$  build up spectrum of higher states. *coupled-channel analysis*
- Enable us to test and verify models of hadronic physics and identify degrees of freedom
- Future prospects. . .

**Dominant cost:**  $m_\pi \longrightarrow 0$

- Straightforward to estimate for **moments of structure functions**
- “Synthetic data” assuming “CSSM” chiral expansion.





### Physics Roadmap at Jefferson Laboratory

